# The collisional formation of massive stars

With Olaf Davis, Marc Freitag & Ian Bonnell

Model of Bonnell, Bate & Zinnecker 1998 (BBZ) proposed because: A) theoretical difficulties with accretion due to p rad (Wolfire & Casinelli 1989; but see Edgar & Clarke 2003, Sonnhalter & Yorke 2002)B) lack of observed discs in massive protostellar cores but see Cesaroni et al 2006 (PPV) et seq.

### Required densities for stellar collisions on Myr timescale is high (~ 10^8 /pc^3)

Stolte et al 2006

Cluster	$M_{ m total}$ $(M_{\odot})$	Extent (pc)	r <sub>core</sub> (pc)	$ ho_{ m core} \ (M_\odot \ { m pc}^{-3})$	Age (Myr)
Arches	104	1 (?)	0.2	$3 \times 10^{5}$	2-3
NGC 3603 YC	$>7 \times 10^{3}$	4.4	0.2	105	1-3
R136	$2 \times 10^{4}$	4.7	0.02	$5 \times 10^{4}$	1-5
Orion	10 <sup>3</sup>	3	0.2	$4 \times 10^{4}$	0.3-1
Antennae starbursts	$10^{4} - 10^{6}$	1-10	?	10 <sup>3</sup>	1-20
Milky Way GCs	10 <sup>4</sup> -few 10 <sup>5</sup>	Few pc	≈1	$10^{2}-10^{6}$	10 Gyr

...collisions unimportant for observed clusters in their current states....

#### Elements of (BBZ) model:

(spherical, non-rotating gas, all gas onto stars)

Adiabatic accretion onto star cluster results in shrinkage:

 $r \propto M^{-3}$ 

Can understand in terms of preservation of adiabatic invariant (e.g. angular momentum)

Or revirialisation after accretion of das with

#### Model scalings

R ~ M^-3

Don't need to accrete much mass to ensure density rises

DRAMATICALLY

• N ~ M^9

• ρ ~ M^ 10

 $V \sim M^{2}$ 

•

Velocity dispersion increases more mildly (=> encounters

remain in grav focused regime)

#### Model limitations

Adiabatic regime requires Mdot < v^3/G (for v = initial velocity cluster vel. dispersion)

More amply satisfied during collapse (once adiabatic => always adiabatic)

If cluster fed by gas free-falling from outer cluster/reservoir, adiabatic condition implies v for cluster > v for reservoir

#### End state

If adiabatic regime persists`forever' are collisions inevitable if keep on accreting?

Bonnell & Bate 2002 found "no" - needed to boost collision cross-section to get collisions before clusters puff up due to few body effects

=> Clarke & Bonnell (2008) argued that collisions require large N, massive clusters (not ONC!)

#### Follow up with Monte Carlo code (Freitag & Benz 2001,2002)

- Models of spherical clusters in dynamical equilibrium
- Include secular effects: a) accretion

Equal mass stars !

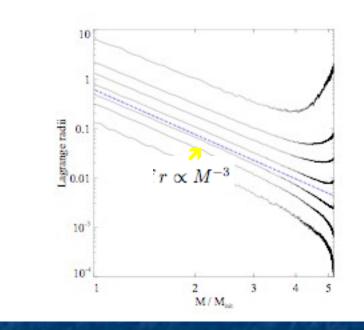
b) two-body relaxation

Estimate collision rate through postprocessing

#### Results

↓

End of adiabatic regime

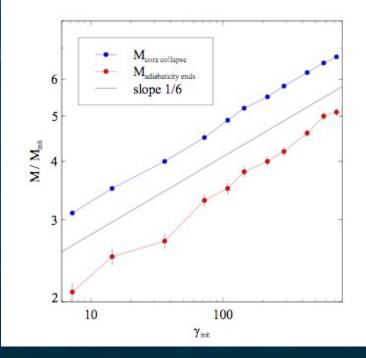




Core collapse!

#### End of adiabatic regime when t\_2r ~ t\_Mdot (= M/Mdot) Ratio of t\_2r/t\_Mdot (initially large) scales as M^-6

 $M_f \propto \gamma_{init}^{rac{1}{6}}.$ Where  $\gamma \equiv rac{t_{2r}}{t_{\dot{M}}}.$ 

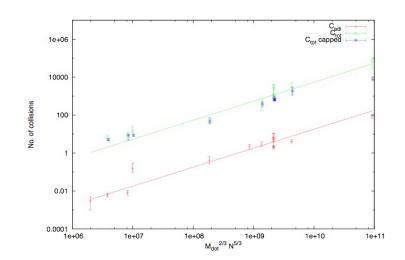




#### Estimate number of collisions in adiabatic phase (collapse homologous) Number of collisions per mass doubling time scales as M^10 during collapse so can relate total collisions to number collisions in first mass doubling time and

 $\gamma \equiv \frac{t_{2r}}{t_{\dot{M}}}$  (since  $M_f \propto \gamma_{init}^{\frac{1}{6}}$ 

find C\_adi proportiona to N^5/3 Mdot^2/3



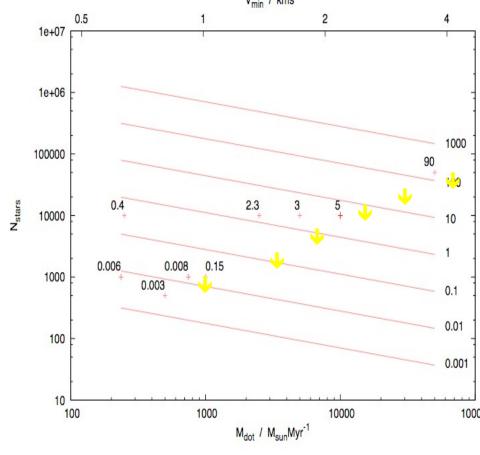
#### Numerical values (red contours) = number of collisions in adiabatic regime initial velocity as function of

V<sub>min</sub> / kms<sup>-1</sup> 0.5 2 1e+07 1e+06 1000 100000 90 10 0.4 2.3 10000 0.006 0.008 0.15 1000 0.003 0.1 0.01 100 0.001 10 100 1000 10000 100000

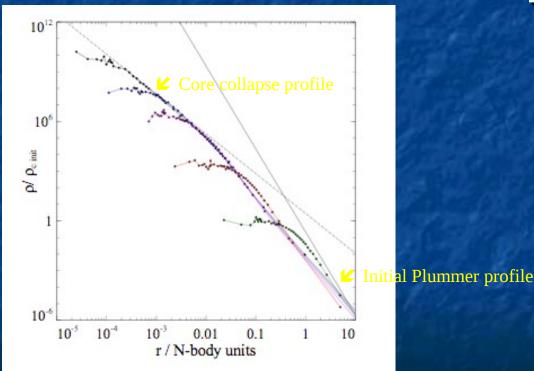
Mdot s.t. enter adiabatic collapse

Have to be below arrows to drive collapse in a Myr **JEED VERY** HIGH N (~ 10^4 -10^5) TO GET EVEN A FEW

COLLISIONS

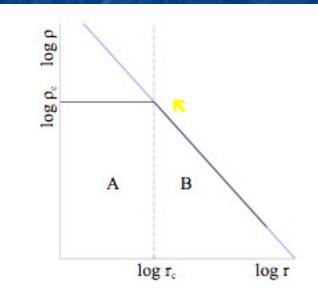


#### But adiabatic regime is not the end of the story - also examine collisions during core collapse



Schematic of density evolution

during core collapse



Collisions predominantly from  $\sim r_c$ 

#### As approach core collapse, system evolves on ~ core relaxation time

Time to core collapse

Evolution of core radius

Evolution of collision rate

 $au = t_{2r}/\eta,$ 

 $r_c \propto ( au \eta)^{5/9}$ 

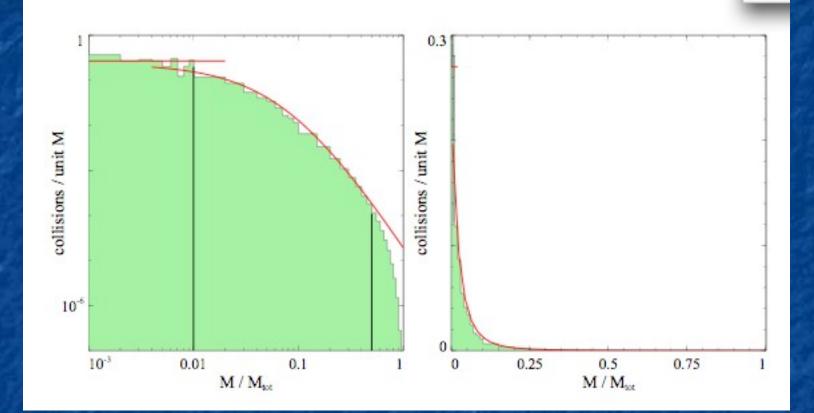
 $\dot{C}_{tot} \propto r_c^{-1.6}.$ 

 $\dot{C} \propto ( au \eta)^{8/9}$ 

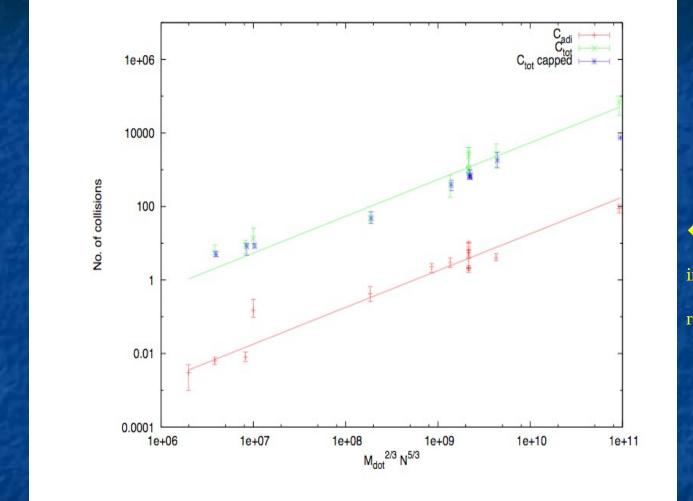
total number of collisions converges

=>

#### Where the collisions occur



Mostly within innermost 10% of mass

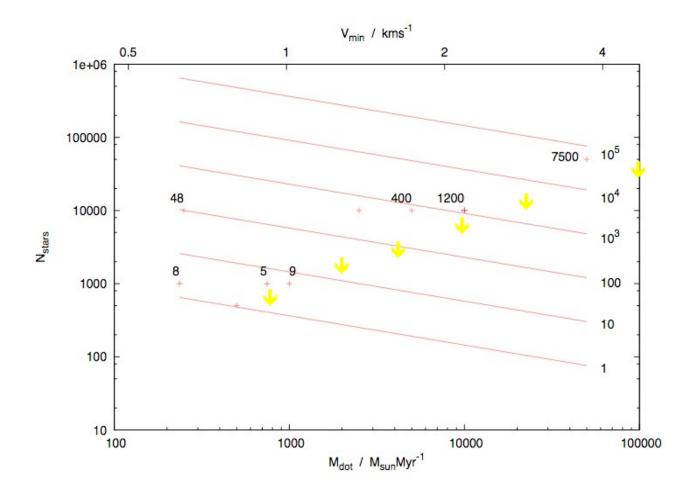




Factor 30!

# Red contours denote number of collisions in core

#### callanca ragima



Interesting number
 of collisions in core
 collapse regime in
 populous clusters

# But would binaries prevent core collapse?

•Three-body binaries only effective in small N clusters

 Primordial binaries? Maybe, but need substantial population of sub-AU binaries

#### Conclusion

•Don't expect many collisions in adiabatic regime even in large N, massive clusters

 But collision yield rises by ~ 30 if core collapse goes to completion- now astronomically interesting in large N, massive systems

Do you get latter phase in realistic case (primordial binaries + mass spectrum)?

#### Adding the realism:

. Include mass spectrum in Monte Carlo (large N) simulations

. Nbody + accretion calculations to include binaries (Nbody6 + "negative stellar winds")