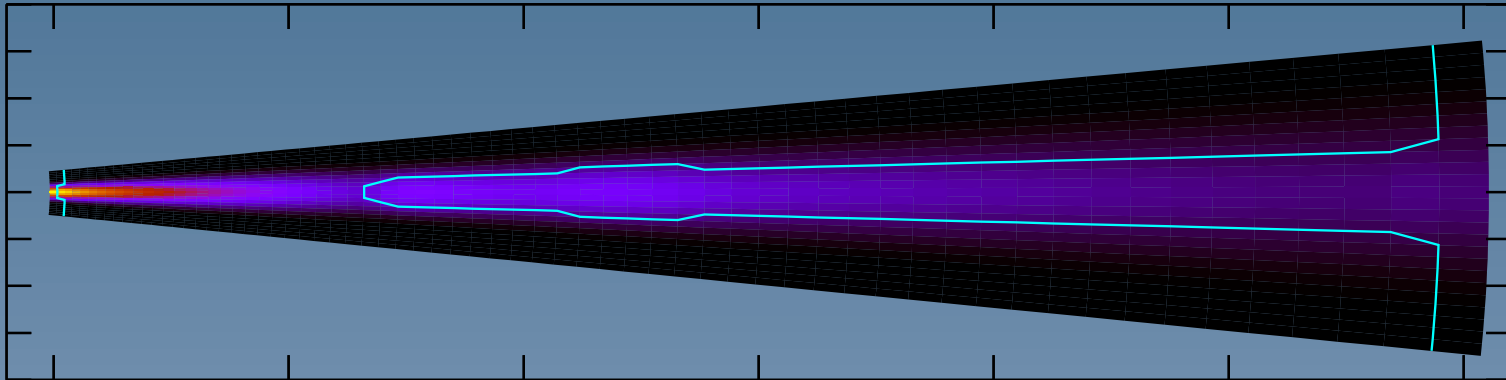


# Layered accretion in protoplanetary disks

(R. Wunsch, M. Różyczka, H. Klahr)



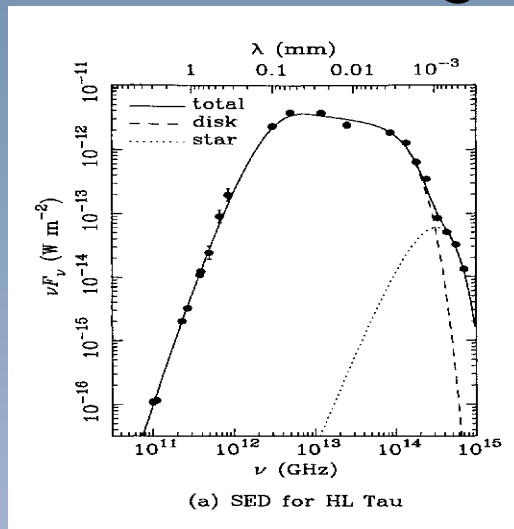
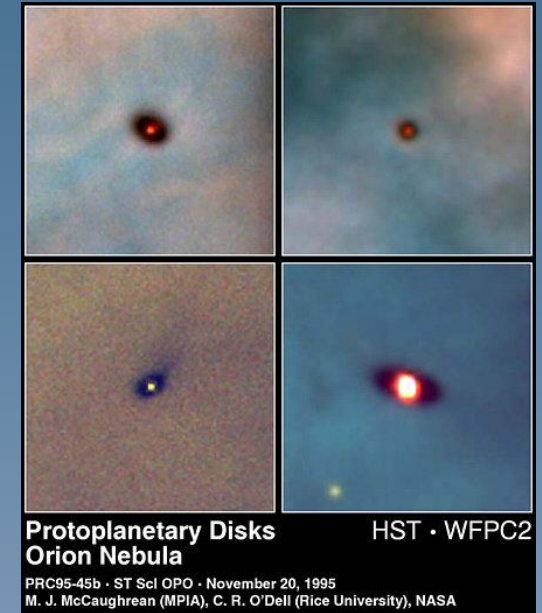
## Outline:

- 1 Protoplanetary disks
- 2 FU Ori outbursts
- 3 Layered-disk model
- 4 Ring instability
- 5 Evolution of the dead zone

# Protoplanetary disks: observational evidence

## Optical/NIR:

- Asymmetric profiles of forbidden emission lines (Edwards et al., 1987)
- Stellar light scattered on dust particles (Beckwith et al., 1989)
- Dark silhouettes in Orion nebulae and other HII regions



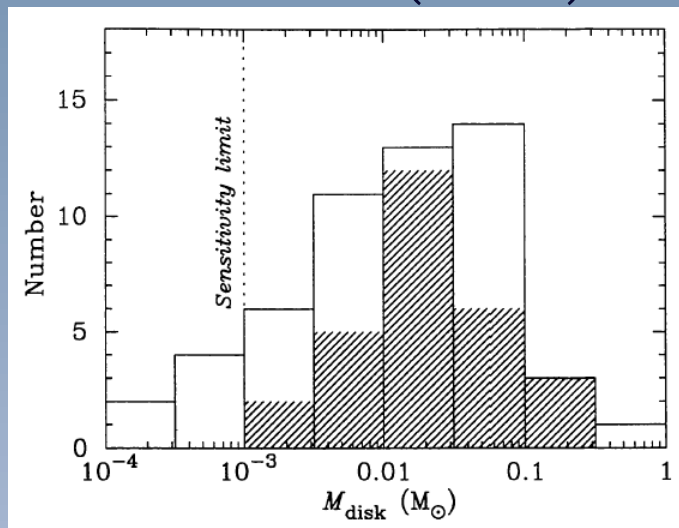
## Millimeter/submillimeter interferometry:

- Broadband emission from the dust (Beckwith et al., 1990)
- CO rotational transitions suggest Keplerian rotational profile

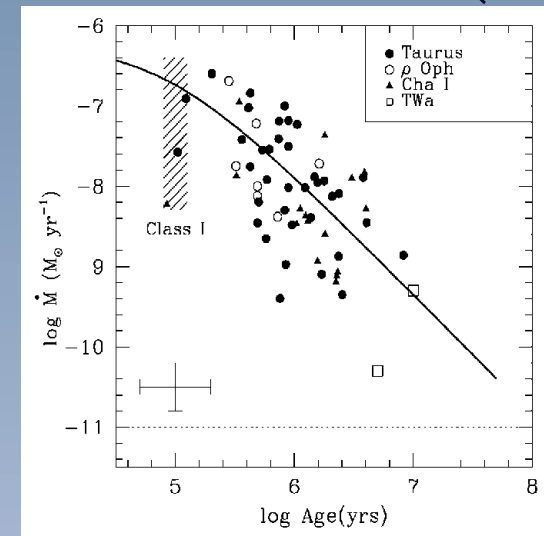
# Protoplanetary disks: observed properties

- Frequency: cca 50% of TTS
- Mass:  $0.01 - 0.1 M_{\odot}$  (from optically thin mm emission)
- Size: 100 - 1000 AU
- Lifetime:  $\sim 10^7$  yr (from ages of TTS)
- Accretion rate:  $10^{-9} - 10^{-6} M_{\odot} \text{yr}^{-1}$   
(from optical/UV excess from inner boundary layer)

Chandler (1998)



Muzerolle et. al (2000)



# Protoplanetary disks: formation and evolution

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1 Collapse of the primordial cloud: in  $\sim 10^4$  yr

2 Formation stage:  $\sim 10^5$  yr

- disk hidden by the stellar gaseous envelope
- energetic mass outflows and jets
- fast evolution due to gravitational instabilities
- finally protostellar envelope dispersed by created pre-MS star

3 Viscous stage:  $\sim 10^7$  yr

- geometrically thin accretion disk, Keplerian rotation
- evolution given by viscosity (probably due to MHD turbulence)
- planets are formed during this stage  $\rightarrow$  protoplanetary disk
- finally disappears (accretion onto star, environmental effects, planet formation)

# Protoplanetary disks: Standard model

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- steady state model:  $\dot{M} = \text{const}$

- mass and angular momentum conservation:

$$\Sigma = \frac{\dot{M}}{3\pi\nu} \left[ 1 - \left( \frac{R_*}{r} \right)^{1/2} \right]$$

- $\alpha$ -viscosity:  $\nu = \alpha c_s H$  (Shakura & Sunayev, 1973)

- energy released by accretion dissipated locally:

$$\frac{9}{8} \Sigma \nu \Omega^2 = \sigma T_e^4$$

- emission is optically thick:  $T_c^4 = \frac{3}{16} \Sigma \kappa T_e^4$  (Hubeny, 1990)

- disk is geometrically thin:  $\Omega = \left( \frac{GM}{R} \right)^{1/2}$

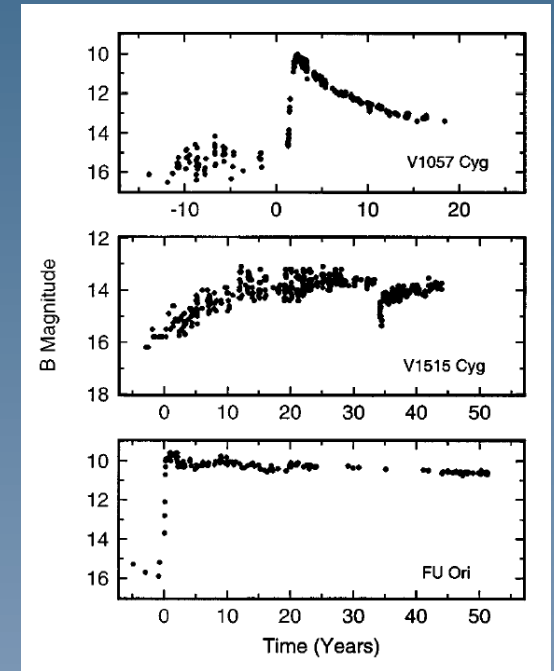
- vertical hydrostatic equilibrium:  $H = \frac{c_s}{\Omega}$

## Solution:

Radial profiles of all interesting quantities ( $\Sigma$ ,  $T_e$ ,  $T_c$ ,  $H$ , ...) are power-laws with exponents dependent on the opacity  $\kappa = \kappa(\rho, T)$ .  $\dot{M}$  and  $\alpha$  are the free parameters.

# FU Ori outbursts

- young stars, large increases in optical brightness  $\sim 4$  mag
- decay timescales: 10-100 yr
- reflection nebulae, heavily extinguished, large IR excess from circumstellar dust
- AD models:  $\dot{M}$  increases from  $10^{-7}$  to  $10^{-4} M_{\odot} \text{ yr}^{-1}$ ,  $\sim 0.01 M_{\odot}$  accreted during one outburst



FU Ori objects (Hartmann & Kenyon, 1996)

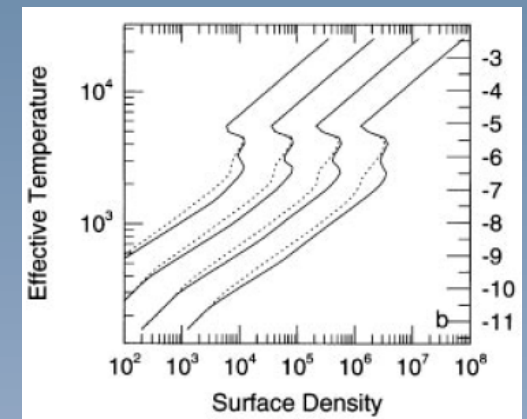
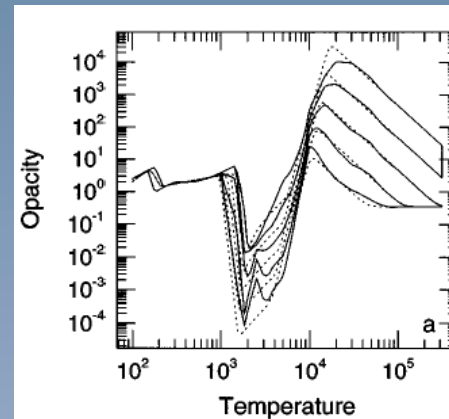
Object	Outburst	$t(\text{Rise})$	$t(\text{Decay})$	$d(\text{kpc})$	$L/L_{\odot}$	CO flow	Jet/HH
FU Ori	1937	$\sim 1$ yr	$\sim 100$ yr	0.5	500	no	no
V1057 Cyg	1970	$\sim 1$ yr	$\sim 10$ yr	0.6	800-250	yes	no
V1515 Cyg	1950s	$\sim 20$ yr	$\sim 30$ yr	1.0	200	no	no
V1735 Cyg	$\sim 1957-65$	$< 8$ yr	$> 20$ yr	0.9	$> 75$	yes	no
V346 Nor	$\geq 1984$	$< 5$ yr	$> 5$ yr	0.7	?	yes	yes
BBW 76	$< 1930$	?	$\sim 40$ yr	1.7?	?	?	no
Z CMa	?	?	$> 100$ yr	1.1	600	yes	yes
L1551 IRS5	?	?	?	0.15	$\geq 20$	yes	yes
RNO 1B,C	?	?	?	0.8	?	yes?	no

# FU Ori outbursts: models

Models have to explain rise times  $\sim 1$  yr  
 $\Rightarrow$  eruption must involve inner region ( $< 1$  AU) of the disk.

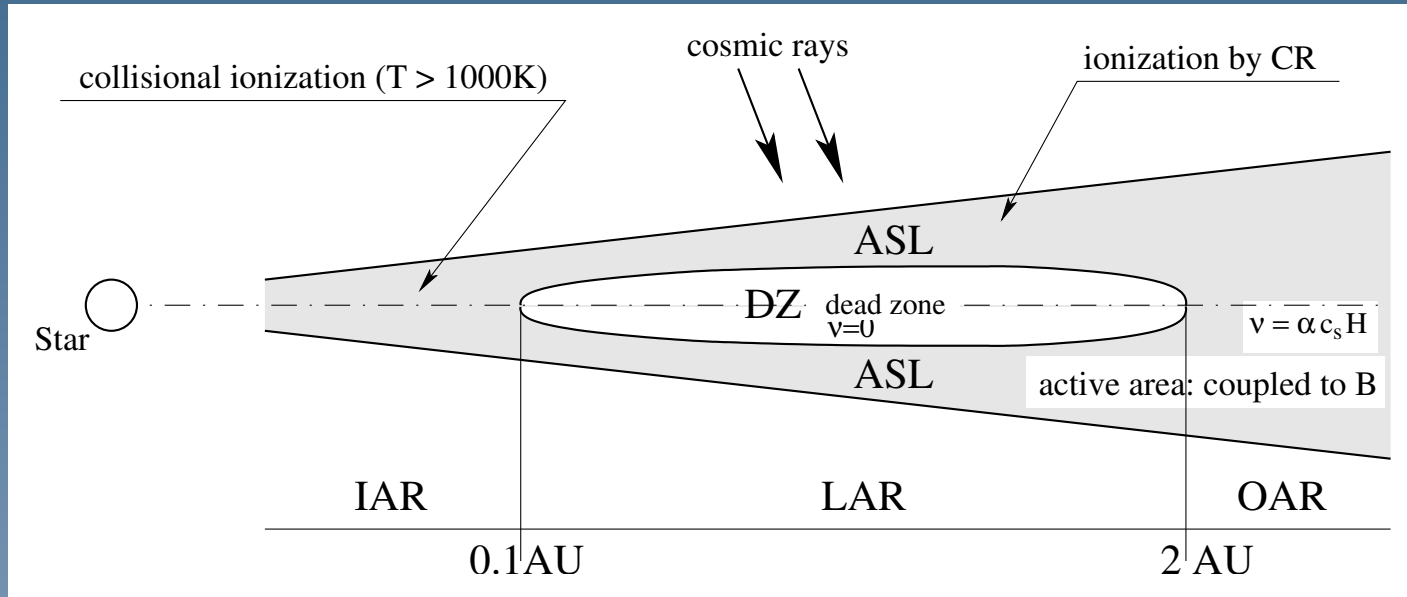
- Triggering of outburst (Bonnel & Bastien, 1992)
  - ▷ disk perturbed by the close passage of the companion star
  - ▷ but no apparent  $v_r$  shifts were observed in the brightest FU Ori objects
- The thermal instability mechanism

- ▷  $F_{\text{vis}}/F_{\text{rad}} \sim \kappa(\rho, T)T^{-3}\Sigma^2\alpha$
- ▷ if  $\kappa$  increases faster than  $T^3$   
 $\Rightarrow$  instability
- ▷ decay timescales too short



- The layered-disk model
  - ▷ accretion of material accumulated in "dead zone"

# Layered-disk: basic idea (Gammie, 1996)



- angular momentum transfer - MRI (Balbus & Hawley, 1991)
- parts of the disk are not ionized enough to be well coupled to the magnetic field
- inner active region (IAR) - collisional ionization
- layered accretion region (LAR) - surface active layers (ASL) ionized by cosmic rays shield the dead zone (DZ) near the mid-plane
- outer active region (OAR) - low surface density, CR are able to ionize whole disk



# Layered-disk: physical processes

- MRI occurs for:  $Re_M \equiv \frac{V_A H}{\eta} > 1$

- Alfvén velocity related to  $\alpha$ -viscosity:  $V_A = \alpha^{1/2} c_s$

- resistivity  $\eta$  related to the ionization degree  $x = n_e/n_H$ :  

$$\eta = 6.5 \times 10^3 x^{-1} \text{cm}^2 \text{s}^{-1}$$

- using  $H = c_s/\Omega$  magnetic Reynolds number:

$$Re_M = 7.4 \times 10^{13} x \alpha^{1/2} \left(\frac{R}{AU}\right)^{3/2} \left(\frac{T}{500K}\right) \left(\frac{M}{M_\odot}\right)^{-1/2}$$

- collisional ionization:  $x = x(\rho, T)$  (Umebayashi, 1983)

$$x \sim \log(\rho), \quad x(T) = \begin{cases} 10^{-16} & \text{for } T \leq 800 \text{ K} \\ 10^{-13} & \text{for } T \sim 900 \text{ K} \\ 10^{-11} & \text{for } T \geq 1000 \text{ K} \end{cases}$$

- CR ionization: stopping depth  $\Sigma_0 \sim 100 \text{ g/cm}^2$

(Umebayashi & Nakano, 1981)

$$x = \left(\frac{\zeta}{\beta n_H}\right)^{1/2} = 1.6 \times 10^{-12} \left(\frac{T}{500K}\right)^{1/4} \left(\frac{\zeta}{10^{-17} \text{s}^{-1}}\right)^{1/2} \left(\frac{n_H}{10^{13} \text{cm}^{-3}}\right)^{-1/2}$$

# Layered disk: analytical solution

- Basic equations for one surface active layer:

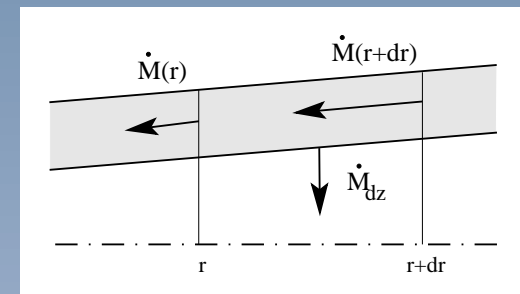
$$\dot{M} = 6\pi r^{1/2} \frac{\partial}{\partial r} (2\Sigma_a \nu r^{1/2}), \quad \frac{9}{4} \Sigma_a \nu \Omega^2 = \sigma T_e^4, \quad \nu = \alpha c_s H$$

$$T_c^4 = \frac{3}{8} \Sigma_a \kappa(\rho, T_c) T_e^4, \quad \Omega = \left( \frac{GM}{R} \right)^{1/2}, \quad H = \frac{c_s}{\Omega}$$

- $\rho_c, T_c, T_e, \dot{M}, H, \dots$  - power-laws
- $\Sigma_a = \text{const} \Rightarrow \dot{M} = \dot{M}(r)$  increasing with  $r$
- accumulation of mass in DZ:

$$\dot{\Sigma}_{\text{DZ}} = \frac{1}{2\pi r} \frac{\partial \dot{M}}{\partial r}$$

- accretion cannot be steady - when  $\Sigma_{\text{DZ}}$  is high enough, DZ mass is accreted in an outburst like event  $\Rightarrow$  FU Orionis



# Numerical model

- RHD code TRAMP: (Klahr et al., 1999)

(Three-dimensional RAdiation-hydrodynamical Modeling Project)

- solves set of Navier-Stokes equations:

$$\frac{\partial \rho}{\partial t} + \nabla \cdot (\rho \mathbf{v}) = 0$$

$$\frac{\partial \rho \mathbf{v}}{\partial t} + (\mathbf{v} \cdot \nabla) \rho \mathbf{v} = -\nabla P - \rho \nabla \Phi + \nabla \cdot \mathbf{T} \quad P = \frac{kT}{\mu m_H} \rho$$

$$c_v \rho \left[ \frac{\partial T}{\partial t} + (\mathbf{v} \cdot \nabla) T \right] = -P \nabla \cdot \mathbf{v} + \mathbf{T} : (\nabla \mathbf{v})$$

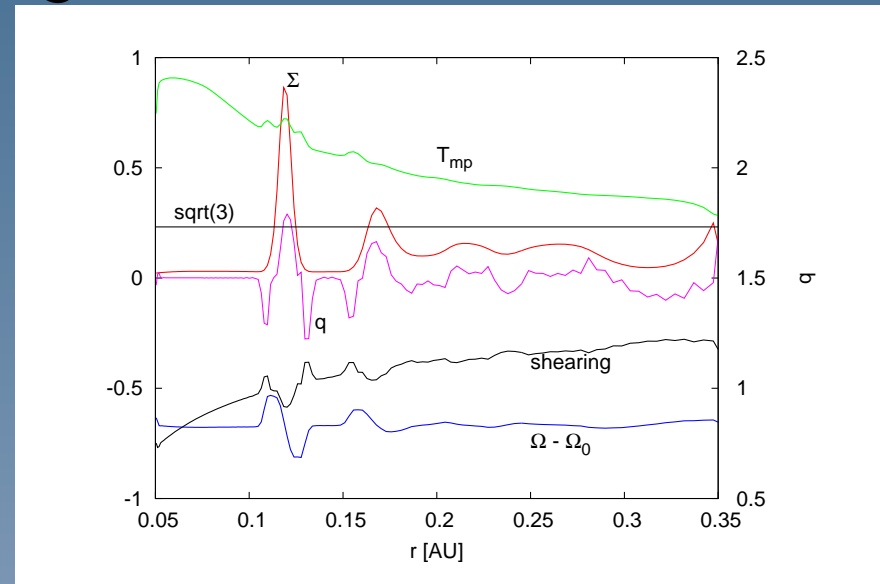
- at the end of time-step: radiation transfer

$$\frac{\partial E_r}{\partial t} = -\nabla \cdot \mathbf{F} \quad , \quad \text{where} \quad E_r = aT^4 \quad , \quad \mathbf{F} = -\frac{\lambda c}{\rho \kappa} \nabla E_r$$

- 2D axially symmetric in spherical  $(r, \theta)$  coords.
- initial conditions: analytical model (vertically isothermal)
- viscosity:
  - ▷  $\alpha = 0.01$  - *surface layers* ( $\Sigma_a = 100 \text{g cm}^{-2}$ ) and inner region ( $T > 1000\text{K}$ )
  - ▷  $\alpha = 0$  - *elsewhere (dead zone)*

# Ring instability

- dead zone decomposes into rings
- ring instability mechanism:
  - ▷ ring-like perturbation changes  $\Omega$
  - ▷ less shearing at the inner edge of the ring
  - ▷ it increases the difference between the amount of mass flowing to the ring from outwards and amount flowing from the ring inwards
  - ▷ positive feedback  $\rightarrow$  growing of the ring



- rings may work as traps for the dust  
 $\rightarrow$  formation of planets
- rings may decay due to the hydrodynamic instability, if  $q > \sqrt{3}$  ( $\Omega \sim r^{-q}$ ) (Papaloizou & Pringle, 1985)

# Ring instability - analytical description

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- thickness of the dead zone important for  $\dot{M}$  (Huré, 2002)  
(vertical gravity determines active layer structure)

- dimensionless disk thickness:  $\delta = \frac{H_a + H_{DZ}}{H_a}$

- rotational velocity corrected to mid-plane pressure:

$$\Omega^2 = \Omega_0^2 + \frac{1}{\rho_m} \frac{\partial P_m}{\partial r}$$

- ring-like perturbation of  $\delta$ :

$$\delta = \delta_0 + \delta_k \cos(kR), \text{ where } R = r - r_0$$

$\delta_0$  - unperturbed disk thickness,  $r_0$  - position of the ring,

$k$  - wavenumber

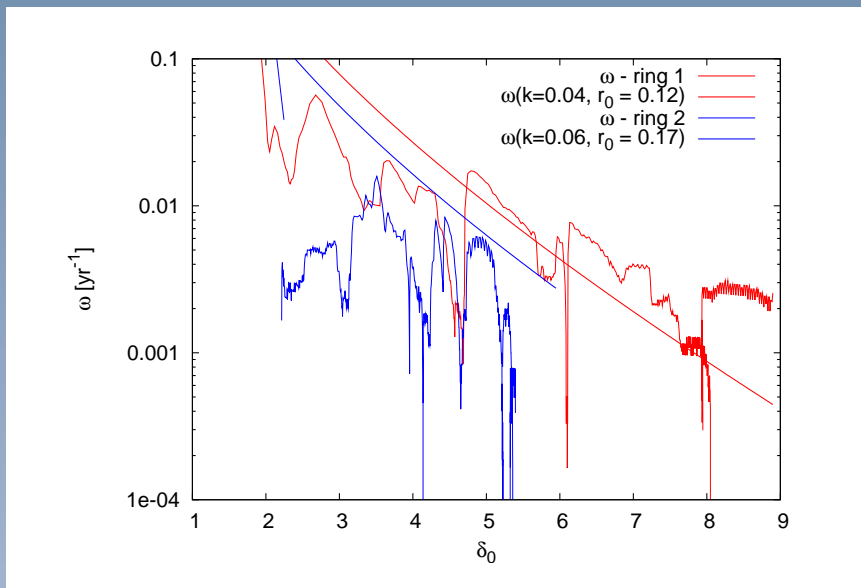
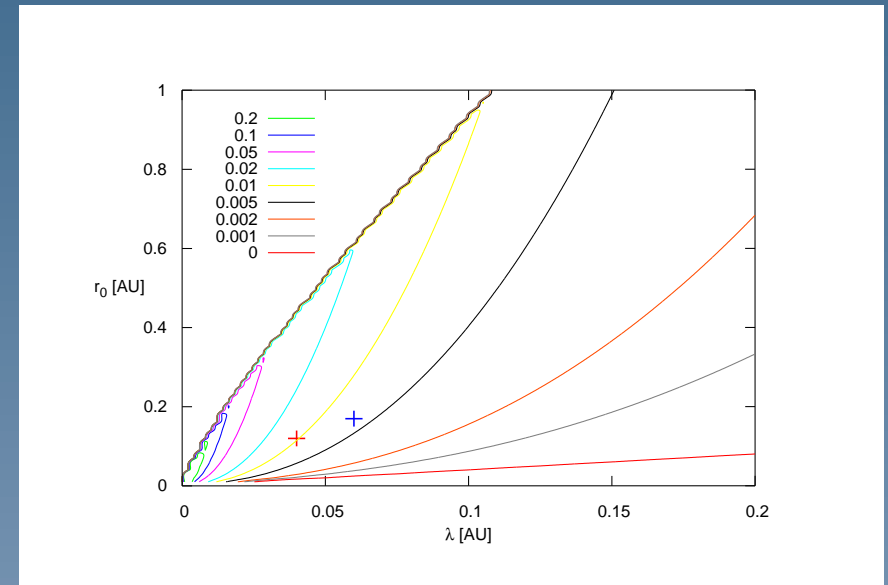
- inserting into equations of layered disk, linearization

- $\rightarrow$  perturbation growth given by the equation:

$$\dot{\delta}_k = \omega(r_0, \delta_0, k) \delta_k$$

# Ring instability - dispersion relation

- $\omega$  diverges for  $k \rightarrow \infty$  ( $\lambda \rightarrow 0$ ) - radial extent of the ring given just by thermal motion of particles
- rings in numerical model several times larger than expected



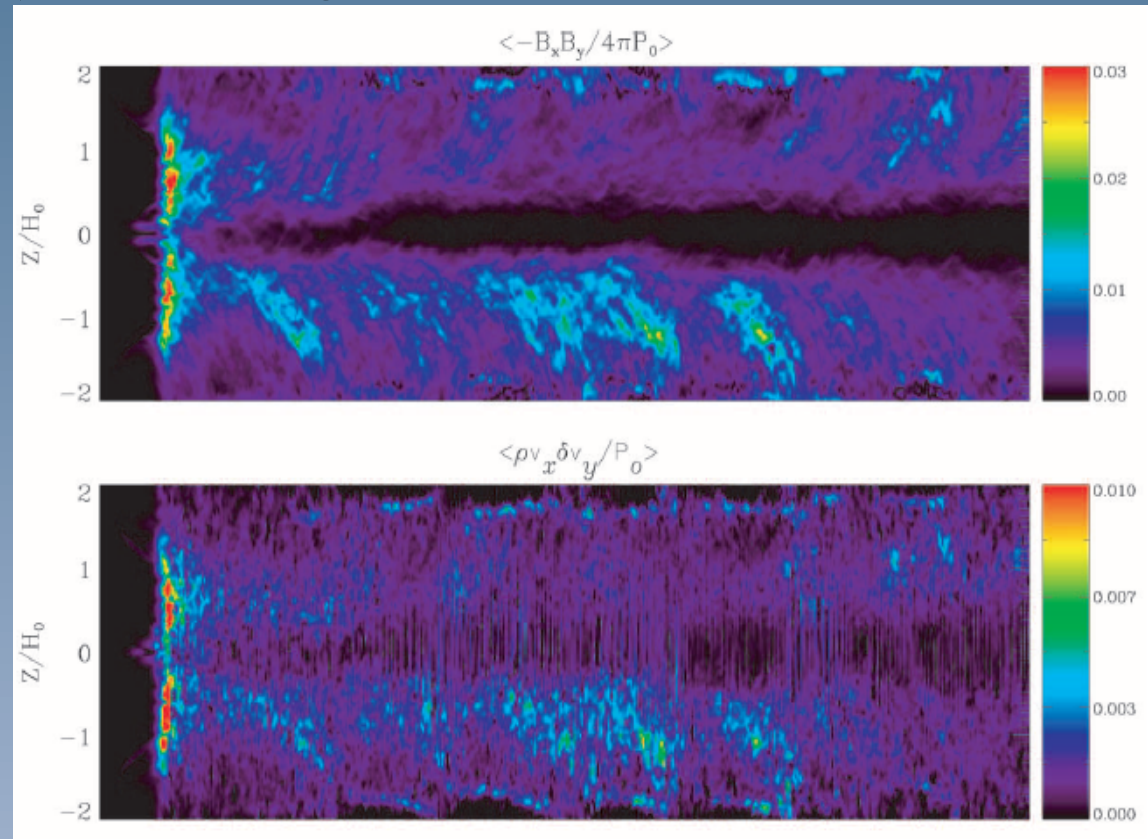
- growth rate of rings and the trend of dependence on  $\delta_0$  roughly in agreement with analytical model

# Layered-disks: MRI simulations

(Fleming & Stone, 2003)

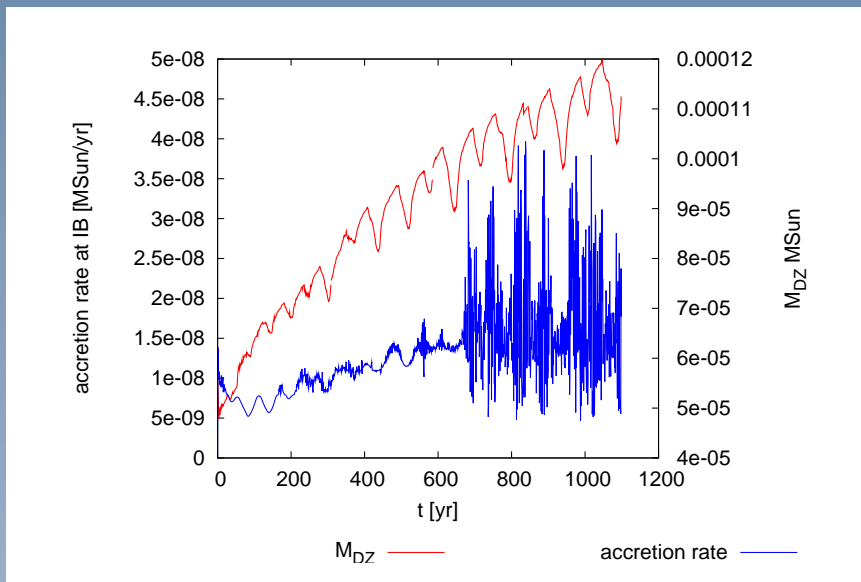
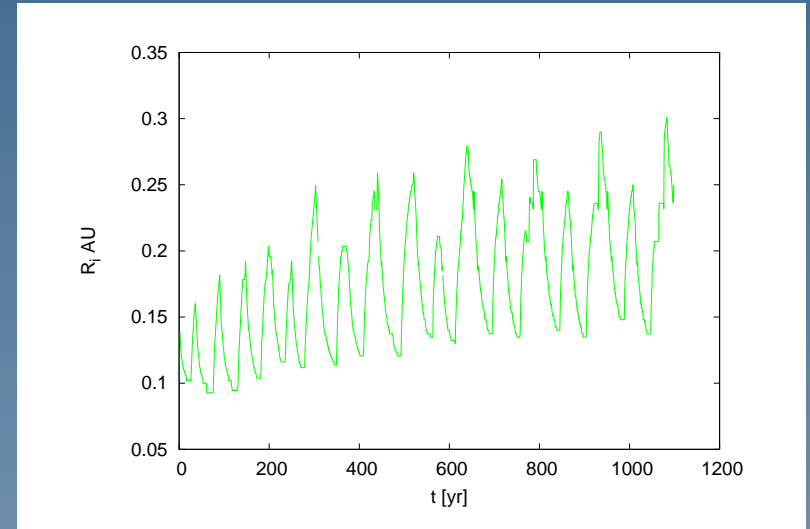
- indications for a small viscosity in the dead zone
- non-ideal MHD ( $\eta \neq 0$ ), shearing box, isothermal EOS

- Maxwell stress vanishes in DZ - MHD turbulence decays
- Reynolds stress non-zero in DZ - HD turbulence survives due to perturbations from active layers (10% of Maxwell s.)



# Evolution of the dead zone

- inner edge of DZ oscillates
- timescale  $\sim 10 - 100$ yr
- DZ continuously depleted from inner edge



- small changes in  $\dot{M}$  (amplitude  $0.5 - 1 \times 10^{-8} M_{\odot} \text{yr}^{-1}$ )
- mass stored in the dead zone grows



# Conclusions

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- the dead zone with zero viscosity is a subject of the ring instability
- in the case of small viscosity in DZ the oscillations of the inner edge occur
- mass is continuously removed from the inner part of the dead zone - it cannot serve as a mass reservoir for the FU Ori type outburst

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