

Planet formation in layered disks

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1. Protoplanetary disks

(observational evidence, properties, formation)

2. Planet formation

(Core accretion vs. Gravitational instability, dust sedimentation, grain growth, accretion of gas, eventual migration, dust-gas dynamics - drag force)

3. Layered disk

(basic model, physical processes, properties of layered accretion, numerical model)

4. Ring instability

(2D simulations, mechanism, analytical description, effect of irradiation)

5. Layered disk evolution - minioutbursts

(residual viscosity in DZ, rings vs. minioutbursts, analytical description, mechanism of mini-outburst, stationary states)

6. Decay of the ring into vortices

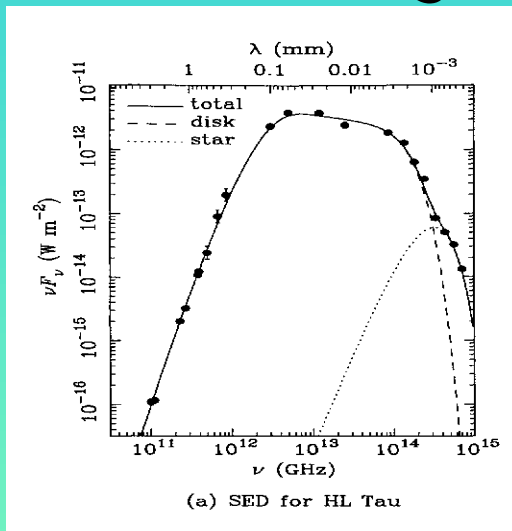
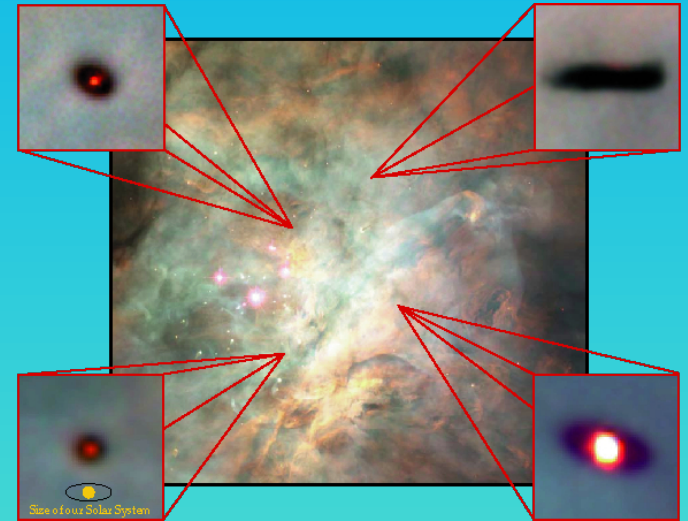
(3D simulations of the Rosby wave instability)

HARDY

Protoplanetary disks: observational evidence

Optical/NIR:

- Asymmetric profiles of forbidden emission lines (Edwards et al., 1987)
- Stellar light scattered on dust particles (Beckwith et al., 1989)
- Dark silhouettes in Orion nebulae and other HII regions



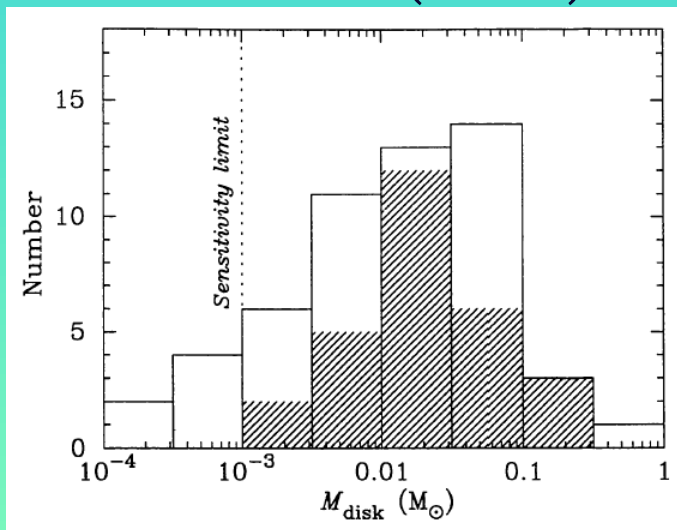
Millimeter/submillimeter, FIR:

- Broadband emission from the dust (Beckwith et al., 1990)
- CO rotational transitions suggest Keplerian rotational profile

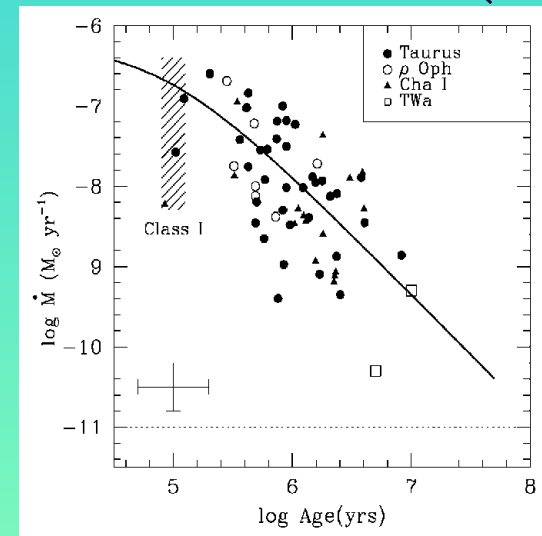
Protoplanetary disks: observed properties

- Frequency: cca 50% of TTS
- Mass: $0.01 - 0.1 M_{\odot}$ (from optically thin mm emission)
- Size: 100 - 1000 AU
- Lifetime: $\sim 10^7$ yr (from ages of TTS)
- Accretion rate: $10^{-9} - 10^{-6} M_{\odot} \text{yr}^{-1}$
(from optical/UV excess from inner boundary layer)

Chandler (1998)

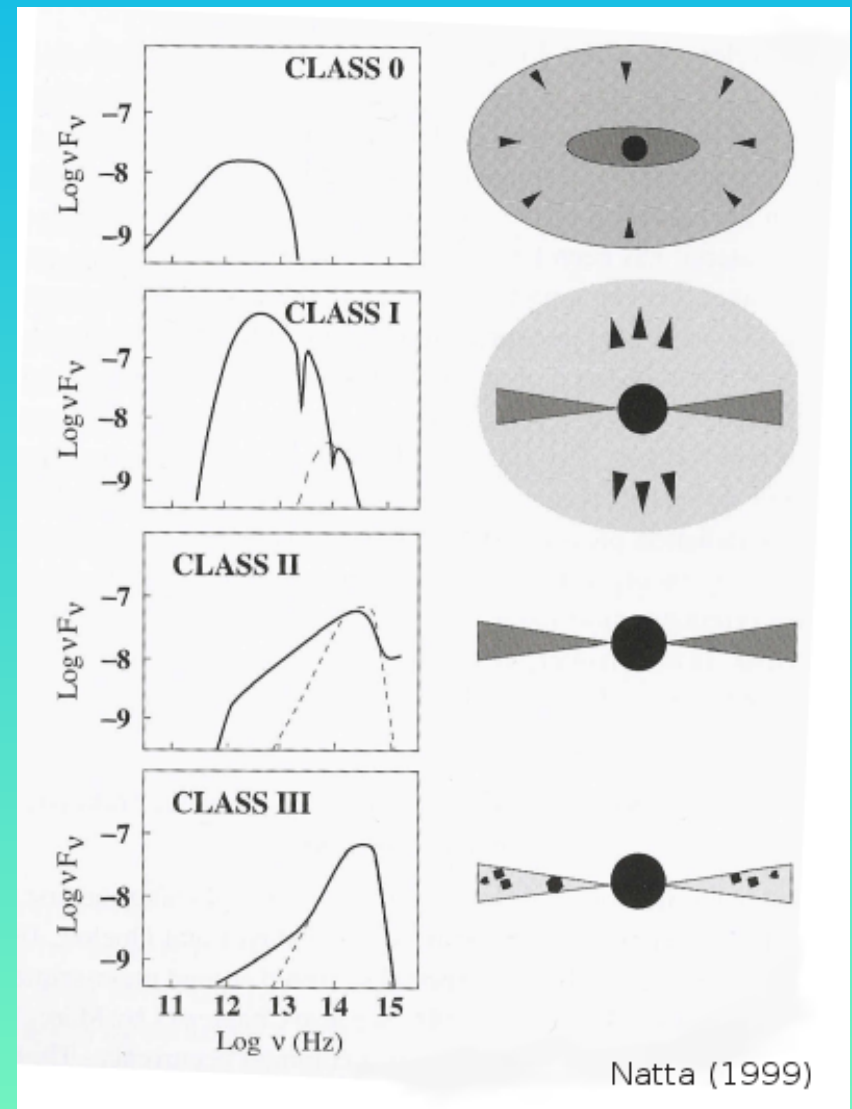


Muzerolle et. al (2000)



Protoplanetary disks: formation and evolution

- **Class 0:**
 - ▶ protostar and disk formed by a collapse of molecular core (10^4 yr)
 - ▶ deeply embedded within an infalling envelope of dust and gas (visible in FIR only)
- **Class I ($\sim 10^5$ yr):**
 - ▶ star accretes matter through the disk, bipolar outflow
 - ▶ still embedded \rightarrow strong absorption features
- **Class II ($\sim 10^7$ yr):**
 - ▶ surrounding material largely dissipated, infall of matter terminated
 - ▶ protostar & optically thick disk
 - ▶ planets are supposed to form
- **Class III:**
 - ▶ gaseous disk almost dispersed (accretion, photo-evaporation)
 - ▶ debris disk: secondary dusty disk formed by solid body collisions



Planet formation: GI or CA?

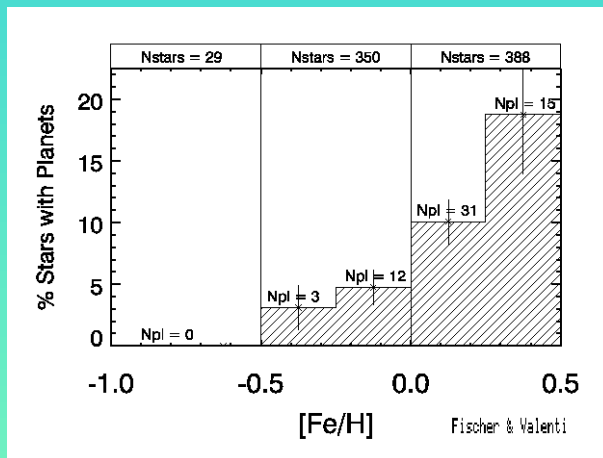
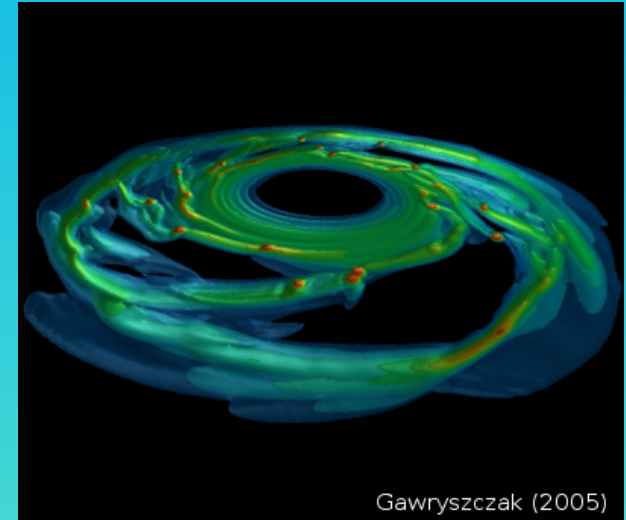
- two competitive hypotheses of giant gas planets formation: **Gravitational instability** (Cameron, 1978)

- massive and cold disk, $Q \equiv \frac{c_s \kappa}{\pi G \sigma} < 1.5$

- fast $\sim 10^3$ yr

- ruled out in mid-1980s - models suggested $M_{\text{Jup,core}} \sim 15 - 30 M_{\oplus}$;

today: $M_{\text{Jup,core}} \sim 6 M_{\oplus}$

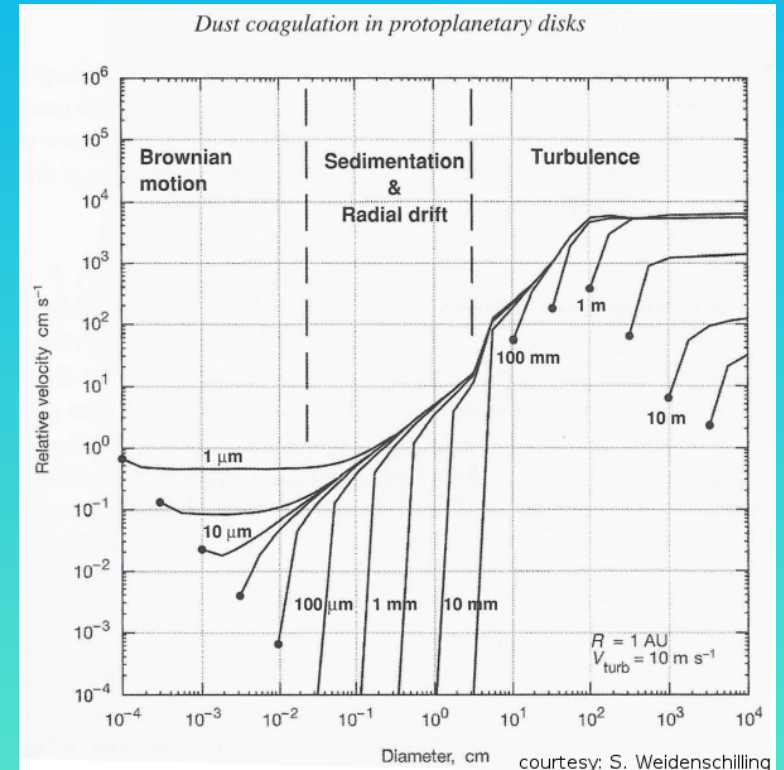


Core accretion (Safronov, 1969)

- collisional accumulation of dust
→ solid core $\sim 10 M_{\oplus}$
- disk gas accretion, ev. migration
- $P(\text{planet})$ depends on Z
- $t_{\text{pl.form.}}$ comparable to disk life-time

Dust coagulation

- $0.1 \mu\text{m}$ grains \rightarrow 100 km bodies
(12 orders of magnitude in size)
 - ▷ complex dust-gas interaction (critical param. $v_{\text{rel}}, \rho_{\text{dust}}$)
 - ▷ **collision:** sticking, bouncing, restructuring, erosion, destruction
 - ▷ Brownian motion, turbulence, sedimentation into dust layer, drag force (radial drift, trapping), gravitational instability of dust layers, ...

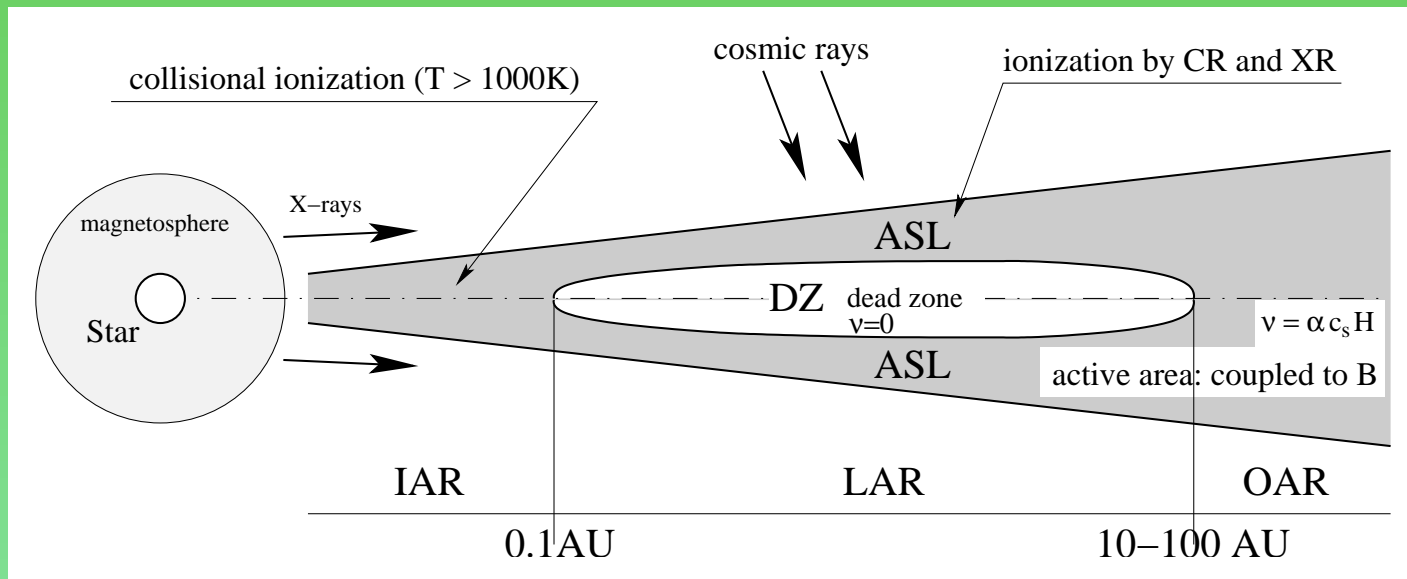


Drag force

$$\left. \begin{aligned} \frac{D\mathbf{v}_{\text{gas}}}{Dt} &= -\nabla P + \rho\nabla\Phi \\ \frac{D\mathbf{v}_{\text{dust}}}{Dt} &= \rho\nabla\Phi \end{aligned} \right\} \frac{D(\mathbf{v}_{\text{dust}} - \mathbf{v}_{\text{gas}})}{Dt} = \nabla P$$

Dust tends to climb up the pressure gradient!
 \rightarrow radial drift, trapping in rings, spirals, vortices

Layered-disk: basic idea (Gammie, 1996)



- angular momentum transfer - MRI (Balbus & Hawley, 1991)
- parts of the disk are not ionized enough to be well coupled to the magnetic field
- inner active region (IAR) - collisional ionization
- layered accretion region (LAR) - surface active layers (ASL) ionized by cosmic rays shield the dead zone (DZ) near the mid-plane
- outer active region (OAR) - low surface density, CR are able to ionize whole disk

Layered-disk: physical processes

- MRI occurs for: $Re_M \equiv \frac{V_A H}{\eta} > 1 - 10^2$ (Fromang et al., 2002)

- Alfvén velocity related to α -viscosity: $V_A = \alpha^{1/2} c_s$

- resistivity η related to the ionization degree $x = n_e/n_H$:

$$\eta = 6.5 \times 10^3 x^{-1} \text{cm}^2 \text{s}^{-1}$$

- using $H = c_s/\Omega$, magnetic Reynolds number is:

$$Re_M = 7.4 \times 10^{13} x \alpha^{1/2} \left(\frac{R}{AU}\right)^{3/2} \left(\frac{T}{500K}\right) \left(\frac{M}{M_\odot}\right)^{-1/2}$$

- collisional ionization: $x = x(\rho, T)$ (Umehayashi, 1983)

$$x \sim \log(\rho), \quad x(T) = \begin{cases} 10^{-16} & \text{for } T \leq 800 \text{ K} \\ 10^{-13} & \text{for } T \sim 900 \text{ K} \\ 10^{-11} & \text{for } T \geq 1000 \text{ K} \end{cases}$$

- CR ionization: stopping depth $\Sigma_0 \sim 100 \text{ g/cm}^2$

(Umehayashi & Nakano, 1981)

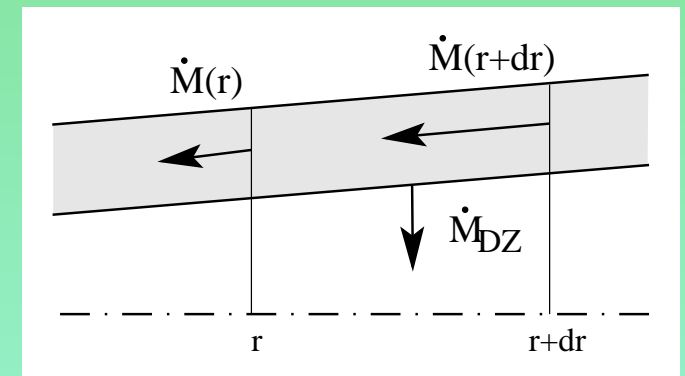
$$x = \left(\frac{\zeta}{\beta n_H}\right)^{1/2} = 1.6 \times 10^{-12} \left(\frac{T}{500K}\right)^{1/4} \left(\frac{\zeta}{10^{-17} \text{s}^{-1}}\right)^{1/2} \left(\frac{n_H}{10^{13} \text{cm}^{-3}}\right)^{-1/2}$$

Properties of the layered accretion

- description of the layered accretion region:

$$\dot{M} = 6\pi r^{1/2} \frac{\partial}{\partial r} (2\Sigma_a \nu r^{1/2}), \quad \nu = \alpha c_{s,i} H_a, \quad \frac{9}{4} \Sigma_a \nu \Omega^2 = \sigma T_e^4$$
$$T_i^4 = \frac{3}{8} \Sigma_a \kappa(\rho_i, T_i) T_e^4, \quad c_{s,i}^2 = H_a (H_a + H_{DZ}) \Omega^2$$

- solution: $\dot{M}(r)$, $T_e(r)$, $T_i(r)$, ... \rightarrow power-laws, exponents dependent on the opacity $\kappa = \kappa(\rho_i, T_i)$
- $\Sigma_a = \text{const} \Rightarrow \dot{M} = \dot{M}(r)$ increasing with r
- accumulation of mass in DZ: $\dot{\Sigma}_{DZ} = \frac{1}{2\pi r} \frac{\partial \dot{M}}{\partial r}$
- accretion cannot be steady - when Σ_{DZ} is high enough, DZ mass is accreted in an outburst like event
 \Rightarrow suggested as mechanism for FU Orionis outbursts (Gammie, 1996)



Numerical model

- based on RHD code TRAMP: (Klahr et al., 1999)
- radiation transfer: flux limited diffusion approximation
- 2D axially symmetric in spherical (r, θ) coords.

- viscosity:

$$\alpha_a = 10^{-2}$$

(surface layers:

$$\Sigma_a = 100 \text{g cm}^{-2}$$

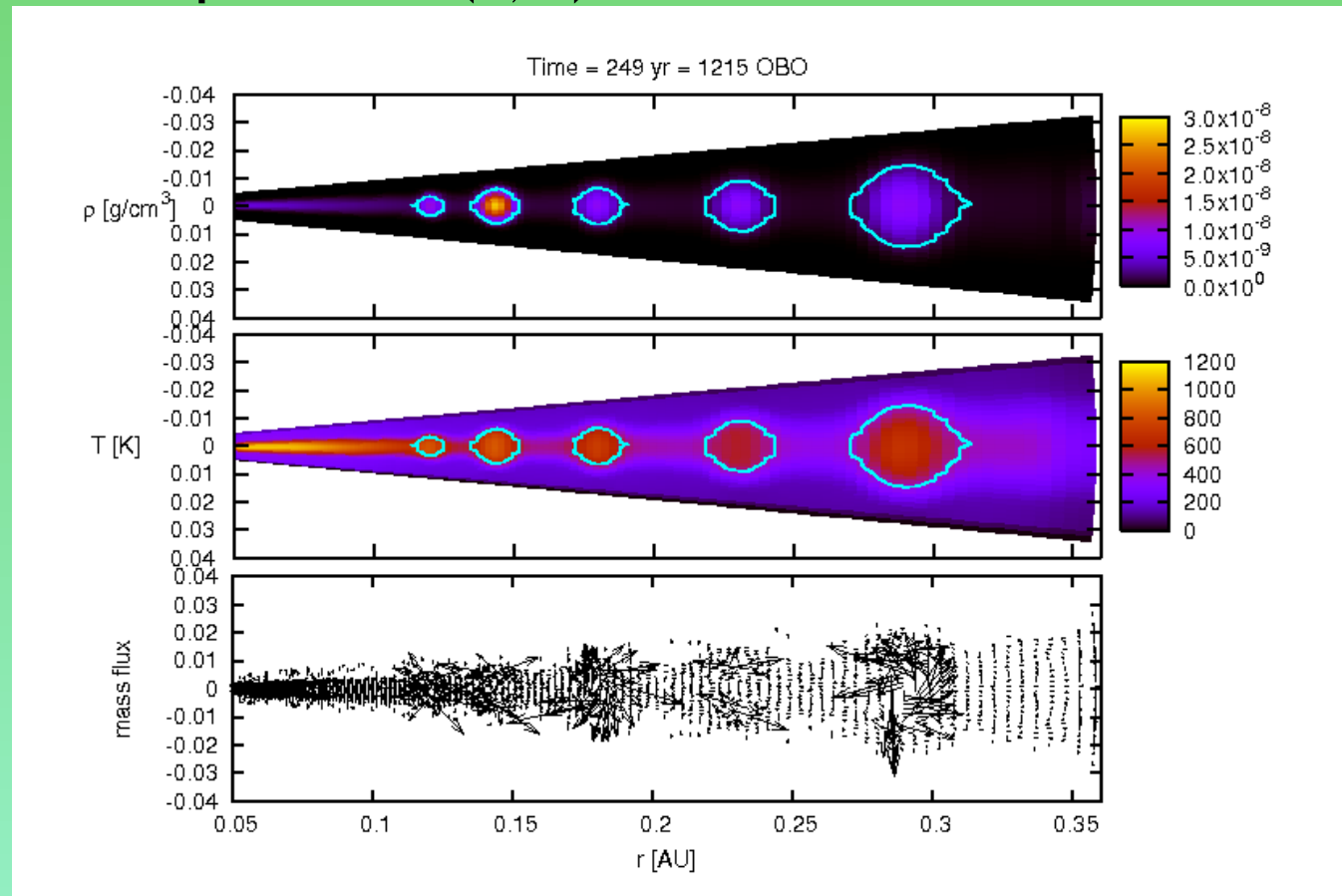
and inner region:

$$T > 1000\text{K}$$

$$\alpha_{DZ} = 0$$

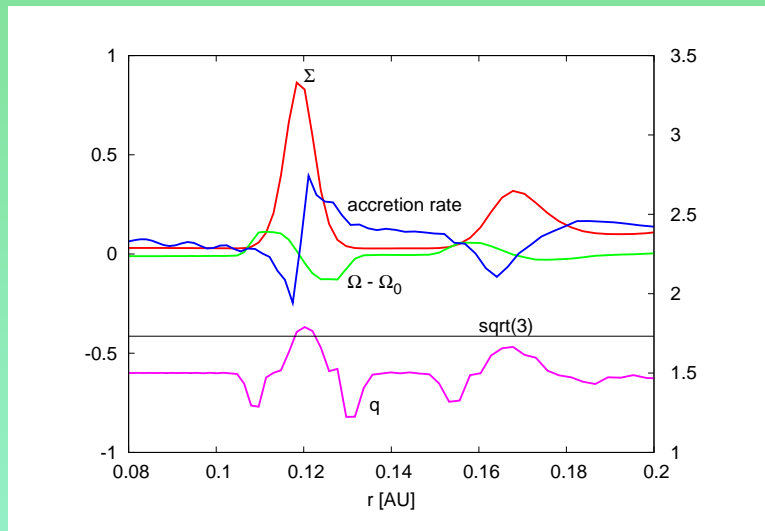
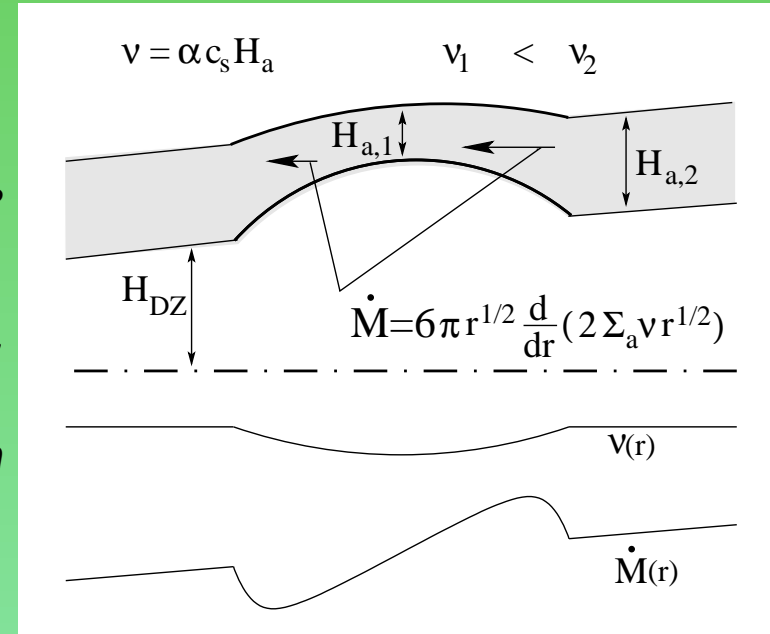
(elsewhere - dead zone)

- rings formed!



Ring instability - mechanism

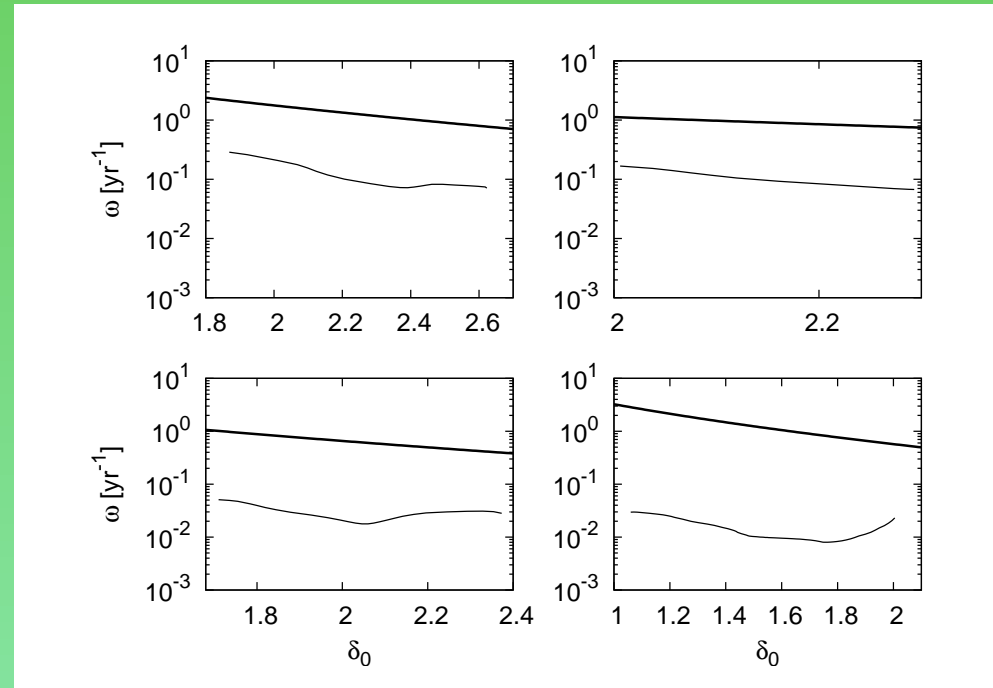
- dead zone decomposes into rings
- ring instability mechanism:
 - ▶ thickness of surface layer H_a depends on the dead zone thickness H_{DZ} (due to different vertical gravity)
 - ▶ H_a is smaller in the ring-like perturbation $\Rightarrow \nu$ is smaller there, too
 - ▶ \dot{M} depends on derivative of $\nu \Rightarrow$ it is smaller in inner edge and higher in outer edge of the ring
 - ▶ enhanced mass accumulation in the ring \Rightarrow positive feedback



- rings may trap the dust \rightarrow higher ρ_{dust} supports dust coagulation
- rings may decay due to the hydrodynamic instability, if $q > \sqrt{3}$ ($\Omega \sim r^{-q}$) (Papaloizou & Pringle, 1985)

Ring instability - anal. description, irradiation

- analytical dispersion relation based on linearized equations was found
- qualitative agreement, numerical ring growth rate is 1-2 orders of magnitude smaller (radiation transport, convective flows make viscosity profiles shallower)

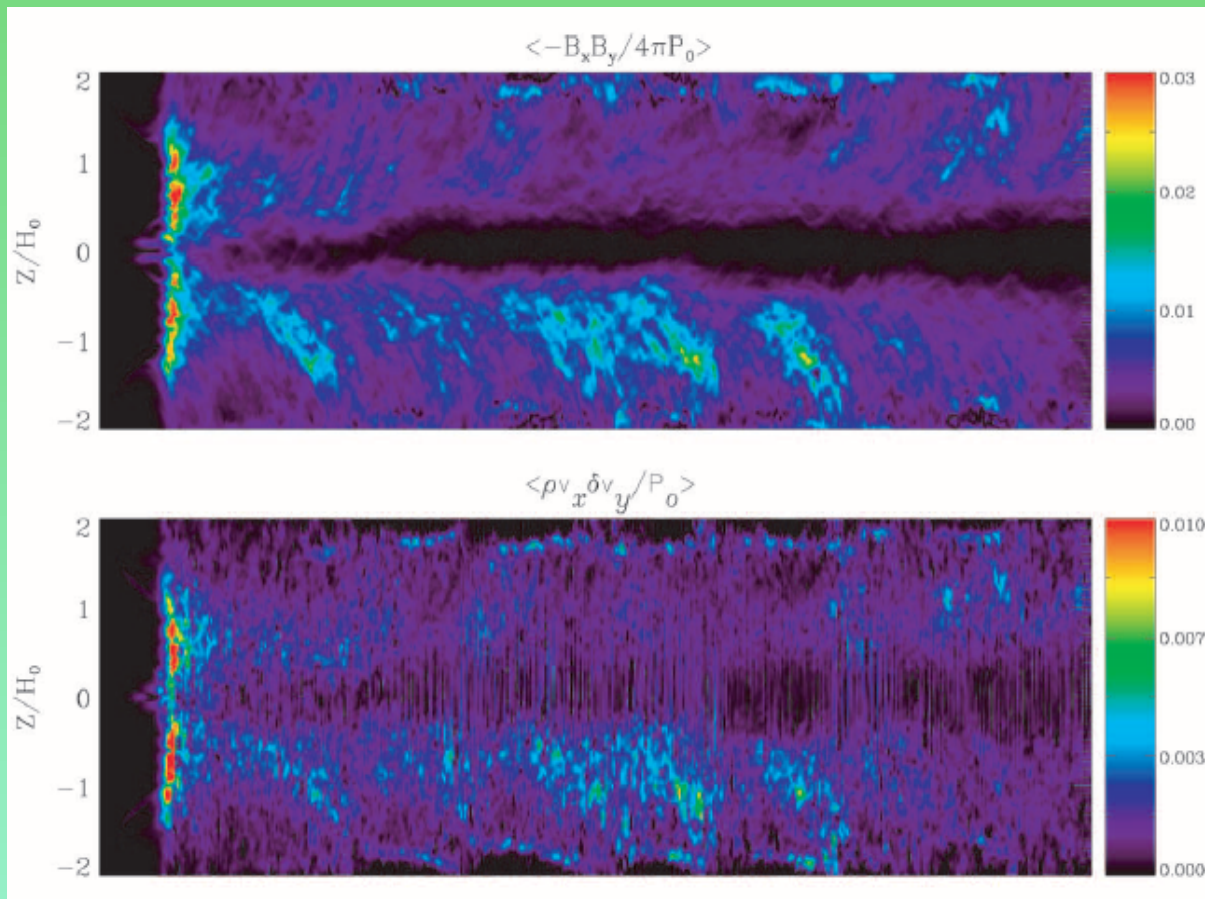


Irradiation

- analytical approximative approach:
$$T_i^4 = \frac{3}{8}\tau T_e^4 + WT_\star^4$$
- vertical structure of the unperturbed disc changes substantially at higher radii ($r > 10$ AU) and for the flaring disc only
- ring inst. can be slowed down, but not suppressed completely

Viscosity in the dead zone - motivation

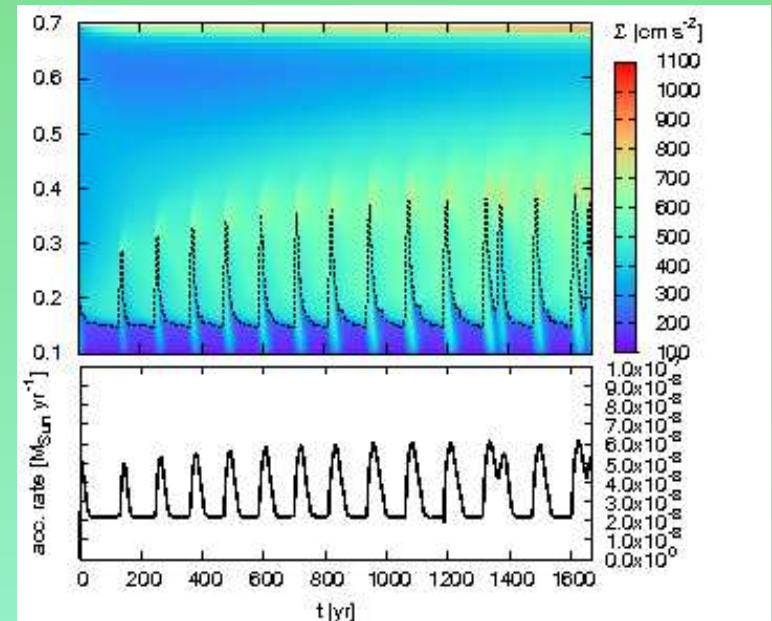
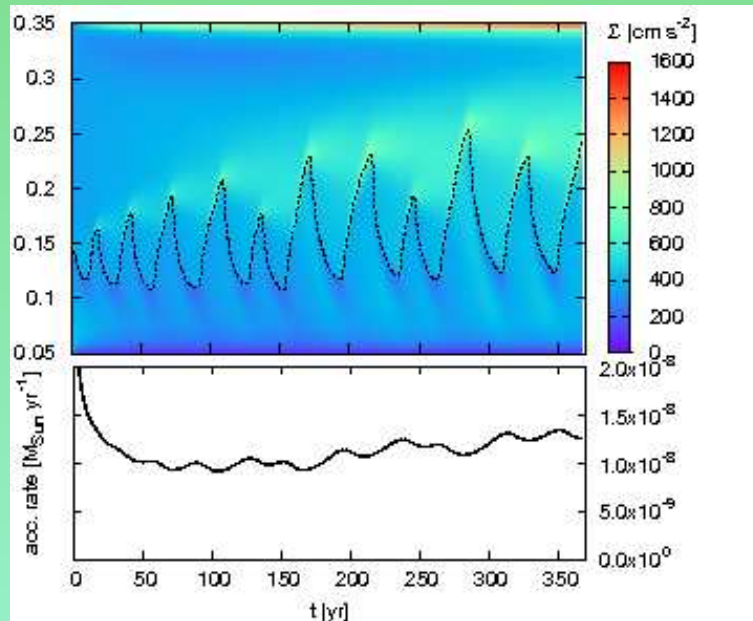
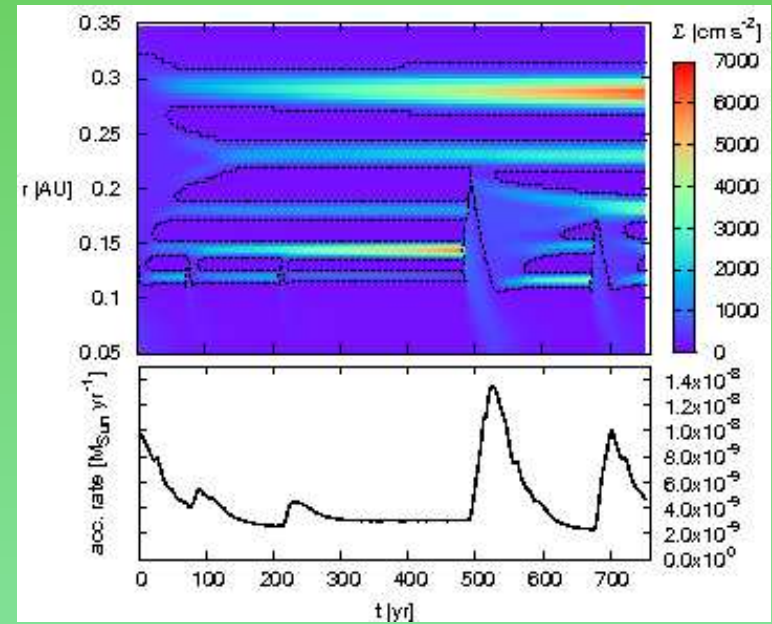
- indications for a small viscosity in the dead zone
- purely hydrodynamic turbulence excited by waves propagating from MRI-active surface layers (Fleming & Stone, 2003)
- viscosity in DZ $\sim 10\%$ viscosity in active parts



- ▷ 3D shearing-box non-ideal MHD simulation
- ▷ time evolution of Maxwell and Reynolds stress tensors along a vertical ray (with a given r and ϕ)
- ▷ Top: Maxwell stress tensor traces MHD turbulence
- ▷ Bottom: Reynolds stress tensor traces pure hydrodynamic turbulence

Results - minioutbursts

- **Right:** $\alpha_a = 0.01$, $\alpha_{DZ} = 0 \Rightarrow$ rings
- **Bottom left:** $\alpha_a = 10^{-2}$, $\alpha_{DZ} = 10^{-3} \Rightarrow$ mini-outbursts, smooth \dot{M}_{IB}
- **Bottom right:** $\alpha_a = 0.02$, $\alpha_{DZ} = 2 \cdot 10^{-3} \Rightarrow$ very short mini-outbursts, \dot{M}_{IB} varies 1 order of mag.



Layered disc with $\alpha_{\text{DZ}} \neq 0$

- analytical description of layered disc with $\alpha_{\text{DZ}} \neq 0$:

$$\dot{M} = 12\pi r^{1/2} \frac{\partial r}{\partial t} (\nu_a \Sigma_a + \nu_{\text{DZ}} \Sigma_{\text{DZ}}), \quad T_m^4 = \frac{3}{8} \kappa T_e^4 \frac{\alpha_a \Sigma_a^2 + \alpha_{\text{DZ}} \Sigma_{\text{DZ}} (\Sigma_{\text{DZ}} + 2\Sigma_a)}{\alpha_a \Sigma_a + \alpha_{\text{DZ}} \Sigma_{\text{DZ}}}$$

$$T_e^4 = \frac{9}{4\sigma} \Omega^2 (\nu_a \Sigma_a + \nu_{\text{DZ}} \Sigma_{\text{DZ}})$$

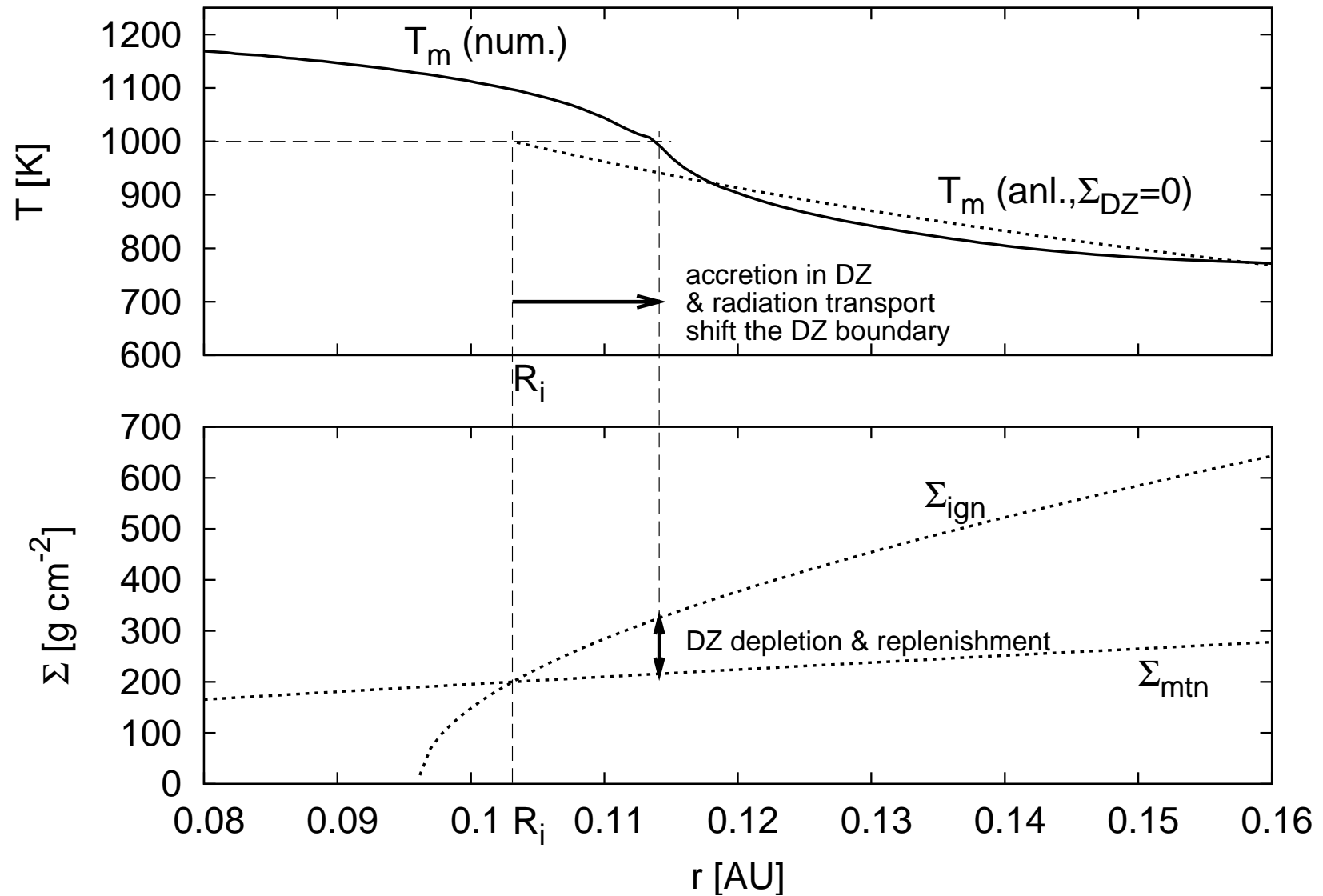
- mid-plane temperature T_m depends on the surface density $\Sigma = 2(\Sigma_a + \Sigma_{\text{DZ}})$ (contrary to LD with $\alpha_{\text{DZ}} = 0$)
- ignition surface density

$$\Sigma_{\text{ign}} = 2 \left(\frac{320\sigma}{27} \frac{\mu m_{\text{H}}}{k_{\text{B}} \Omega \alpha_{\text{DZ}}} T_{\text{lim}}^{5/2} - \frac{\alpha_a - \alpha_{\text{DZ}}}{\alpha_{\text{DZ}}} \Sigma_a^2 \right)^{1/2}$$

- surface density necessary to maintain the disc active

$$\Sigma_{\text{mtn}} = \left(\frac{1280\sigma \mu m_{\text{H}}}{27 k_{\text{B}} \alpha_a \Omega} \right)^{1/2} T_{\text{lim}}^{5/4}$$

Mini-outbursts - mechanism



Stationary states

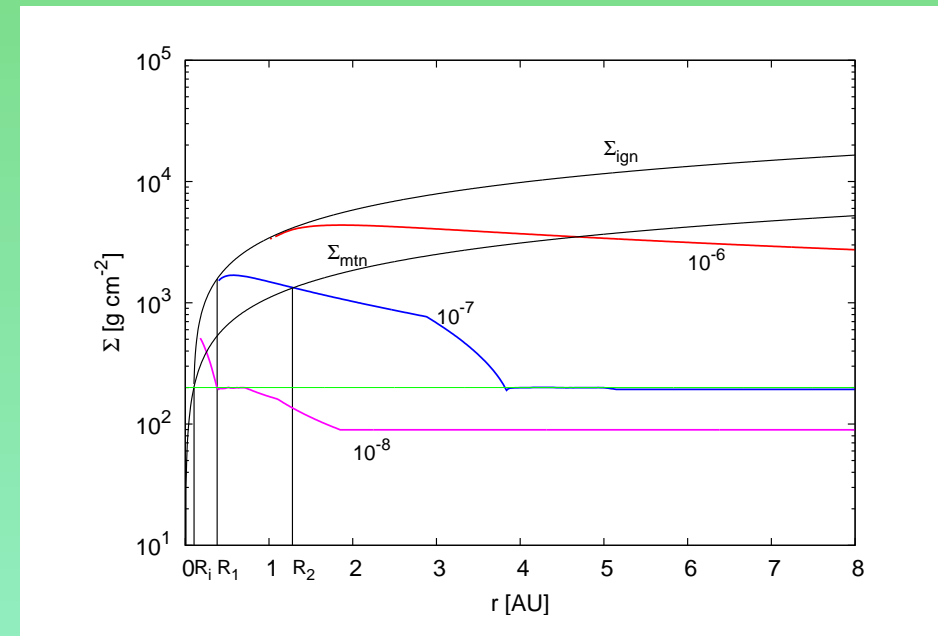
- we search for stationary surface density profiles using 1D code similar to Stepinski (1999) or Armitage et al. (2001)

- numerically solve equation
$$\dot{\Sigma} = \frac{1}{2\pi r} \frac{\partial \dot{M}}{\partial r}$$

- determines T_m from Σ , decides if LD or α D, computes \dot{M} , advects mass between cells

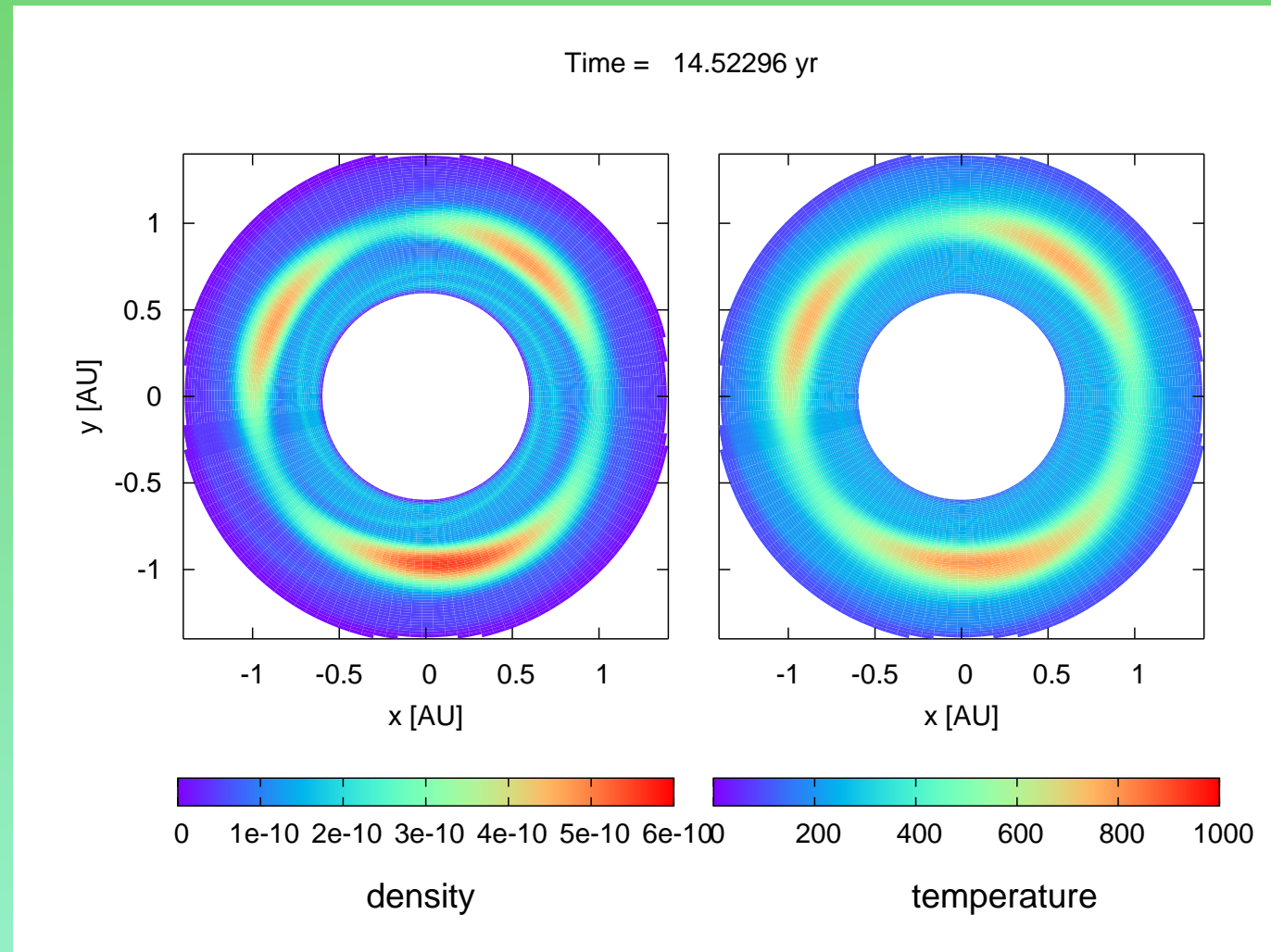
- ▷ stationary solution exist above R_1
- ▷ part of DZ between R_1 and R_2 can be ignited externally
- ▷ it contains mass

$2.0 \times 10^{-7} M_{\odot}$ for $\dot{M} = 10^{-8} M_{\odot} \text{ yr}^{-1}$
 $2.7 \times 10^{-4} M_{\odot}$ for $\dot{M} = 10^{-7} M_{\odot} \text{ yr}^{-1}$
 $0.01 M_{\odot}$ for $\dot{M} = 10^{-6} M_{\odot} \text{ yr}^{-1}$



Decay of the ring into vortices

- 3D simulation of the ring (64x51x128)
- Rossby wave instability (shear, similar to Helmholtz-Kelvin)
 - ▷ *analytical description* (Papaloizou & Pringle, 1985; Lovelace et al., 1999), *2D simulations* (Hawley, 1987; Finn et al., 2000)
 - ▷ *does it work in 3D?*
 - ▷ *for which parameters?* (viscosity, opacity law, resolution)
 - ▷ *combined HD - N-body simulations suggest dust particle concentration* (Johansen, et al., 2007)



Planet formation in vortices - a new idea?

Descartes (1644): Concerning the creation of all of the Planets:

". . . the extremely large space which now contains the vortex of the first heaven was formerly divided into fourteen or more vortices. . . So that since those three vortices which had at their centers those bodies that we now call the Sun, Jupiter, and Saturn were larger than the others; the stars in the centers of the four smaller vortices surrounding Jupiter descended toward Jupiter. . .

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