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Dynamics in the vicinity of the circular trajectory:

cylindrical coordinates: $(R, \psi, z, \dot{R}, \dot{\psi}, \dot{z})$

Assumptions:

1. $\Phi \neq \Phi(t)$ - quasiequilibrium
2. $\Phi = \Phi(\psi)$ - axial symmetry
3. $\Phi(R, z) = \Phi(R, -z)$ - plane symmetry
4. $\Phi(R, z) = \Phi_1(R) + \Phi_2(z)$ - pot. is separable

Integrals of motion:

$$\bar{E} = \frac{1}{2} (\dot{R}^2 + R^2 \dot{\psi}^2 + \dot{z}^2) + \bar{\Phi}(R, z)$$

$$h = R^2 \dot{\psi}$$

Equations of motion:

$$\text{from } \frac{\partial}{\partial t} \left(\frac{\partial L}{\partial \dot{q}} \right) - \frac{\partial L}{\partial q} = 0, \quad L = \frac{1}{2} (\dot{R}^2 + R^2 \dot{\psi}^2 + \dot{z}^2) - \bar{\Phi}(R, z)$$

$$\ddot{R} = R \dot{\psi}^2 - \frac{\partial \bar{\Phi}}{\partial R}$$

$$\frac{d}{dt} (R^2 \dot{\psi}) = - \frac{\partial \bar{\Phi}}{\partial \psi} = 0$$

$$\ddot{z} = - \frac{\partial \bar{\Phi}}{\partial z}$$

Circular orbit:

$$\ddot{R}_0 = 0 \Rightarrow R_0 \dot{\psi}_0^2 = \left(\frac{\partial \bar{\Phi}}{\partial R} \right)_0 = \frac{\theta_0^2}{R_0}$$

$$h_0^2 = R_0^2 \theta_0^2 = R_0^3 \left(\frac{\partial \bar{\Phi}}{\partial R} \right)_0 \Rightarrow h_0 = R_0^2 \dot{\psi}_0$$

Nearly circular orbit:

$$R = R_0 + R_1, \quad R_1 \ll R_0$$

$$\ddot{R} = (R_0 + R_1) \dot{\psi}^2 - \frac{\partial \bar{\Phi}}{\partial R}$$

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$$\ddot{R}_1 = (R_0 + R_1) \dot{v}^2 - \left(\frac{\partial \Phi}{\partial R}\right)_0 - R_1 \left(\frac{\partial^2 \Phi}{\partial R^2}\right)_0$$

$$\ddot{R}_1 = \frac{h^2}{(R_0 + R_1)^3} - \left(\frac{\partial \Phi}{\partial R}\right)_0 - R_1 \left(\frac{\partial^2 \Phi}{\partial R^2}\right)_0$$

$$\ddot{R}_1 = \frac{h^2}{R_0^3} \left(1 - 3 \frac{R_1}{R_0}\right) - \left(\frac{\partial \Phi}{\partial R}\right)_0 - R_1 \left(\frac{\partial^2 \Phi}{\partial R^2}\right)_0 \quad h_0^2 = R_0^3 \left(\frac{\partial \Phi}{\partial R}\right)_0$$

$$\ddot{R}_1 = -3 \frac{h_0^2}{R_0^3} \frac{R_1}{R_0} - R_1 \left(\frac{\partial^2 \Phi}{\partial R^2}\right)_0$$

taking $h = h_0$ - on this circular orbit; other h belongs to different circ. orbit

$$\ddot{R}_1 = - \left[\frac{3}{R_0} \left(\frac{\partial \Phi}{\partial R}\right)_0 + \left(\frac{\partial^2 \Phi}{\partial R^2}\right)_0 \right] R_1$$

$$\boxed{\ddot{R}_1 = -\kappa^2 R_1} \quad \kappa^2 - \text{epicyclic frequency}$$

stability for $\kappa^2 > 0$:

$$\kappa^2 R_0^3 = \left[\frac{\partial}{\partial R} \left(R^3 \frac{\partial \Phi}{\partial R} \right) \right]_0$$

$$\frac{\partial \Phi}{\partial R} \sim R^\alpha, \quad \alpha > -2 \text{ always} \Rightarrow \text{orbit is always stable}$$

Using Oort's consts:

$$F_{R_0} = \left(\frac{\partial \Phi}{\partial R}\right)_0 = \frac{\theta_0^2}{R_0}, \quad \left(\frac{\partial F_R}{\partial R}\right)_0 = \left(\frac{\partial^2 \Phi}{\partial R^2}\right)_0 = -\frac{\theta_0^2}{R_0^2} + 2 \frac{\theta_0}{R_0} \left(\frac{d\theta}{dR}\right)_0$$

$$\kappa^2 = \frac{3}{R_0} \left(\frac{\partial \Phi}{\partial R}\right)_0 + \left(\frac{\partial^2 \Phi}{\partial R^2}\right)_0 = \frac{3}{R_0} \frac{\theta_0^2}{R_0} - \frac{\theta_0^2}{R_0^2} + 2 \frac{\theta_0}{R_0} \left(\frac{d\theta}{dR}\right)_0$$

$$\kappa^2 = 2 \frac{\theta_0}{R_0} \left[\underbrace{\frac{\theta_0}{R_0}}_{(A-B)} + \underbrace{\left(\frac{d\theta}{dR}\right)_0}_{2B} \right] = 4(A-B)B$$

$$\kappa^2 = 1300, \quad \kappa = 36 \text{ km s}^{-1} \text{ kpc}^{-1}, \quad T_{ep} = \frac{2\pi}{\kappa} = 170 \text{ Myr}$$