

1. Modified Boyle's law by Terletsky

$pV = NkT$ - Boyle's law

$pV = NkT - \alpha GM^2 V^{-1/3}$ - correction for grav. energy (Terletsky, 1952)

volume fluctuations of the ideal gas given by:

$$\frac{\overline{(V - \bar{V})^2}}{\bar{V}^2} = - \frac{kT}{V^2 \left(\frac{\partial p}{\partial V}\right)_T}$$

(Statistical Physics Landau & Lifshitz §112 Fluctuations of the fundamental thermodyn. q. eq. 112.7)

$$p = \frac{NkT}{V} - \alpha \frac{GM^2}{V^{1/3}}$$

$$\left(\frac{\partial p}{\partial V}\right)_T = - \frac{NkT}{V^2} + \frac{4\alpha}{3} \frac{GM^2}{V^{4/3}} \quad \leftarrow M = mN$$

$$\frac{\overline{(V - \bar{V})^2}}{\bar{V}^2} = \frac{1}{N} \left[1 - \frac{4}{3} \frac{\alpha G m^2 N}{kT V^{1/3}} \right]^{-1}$$

fluktuations large if $\left[\begin{array}{l} N \text{ is small} \\ \frac{M}{R} = \frac{mN}{V^{1/3}} \approx \frac{3kT}{4\alpha Gm} \end{array} \right.$

2. Hydrostatic equilibrium eq.

(2)

$$-\frac{dp}{dr} = \frac{4\pi G \rho}{r^2} \int_0^r \rho y^2 dy$$

$$+ p = \frac{\rho k T}{m}$$

$$\frac{1}{r^2} \frac{d}{dr} \left(\frac{r^2}{\rho} \frac{dp}{dr} \right) = -4\pi G \rho$$

→ Emden's equation:

$$\frac{1}{r^2} \frac{d}{dr} \left(\frac{r^2}{\rho} \frac{dp}{dr} \right) = - \frac{4\pi G m \rho}{k T}$$

→ in a form of the Boyle's law:

$$N = \frac{4\pi}{m} \int_0^r \rho y^2 dy$$

$$p = \frac{\rho k T}{m}$$

$$pV - NkT = \frac{4\pi}{3} r^3 p - \frac{4\pi k T}{m} \int_0^r \rho y^2 dy =$$

$$= \frac{4\pi}{3} \int_0^r y^3 \frac{dp}{dy} dy$$

$$pV = NkT - \frac{16}{3} \pi^2 G \int_0^r \rho y dy \int_0^y \rho x^2 dx$$

for $\rho = \text{const}$ (which does not satisfy hydrost. equil.) we get:

$$\alpha = \frac{1}{5} \left(\frac{4\pi}{3} \right)^{1/3}$$

3. Transformation of Emden's equation:

(3)

$$r, \rho, T, m \rightarrow \xi, \psi, \lambda, \beta$$

$$r = \beta^{1/2} \lambda^{-1/2} \xi, \quad \rho = \lambda \exp(-\psi), \quad \beta = \frac{kT}{4\pi G m}$$

Boundary cond. at $r=0$:

$$\begin{aligned} \rho = \rho_c & \Rightarrow \lambda = \rho_c \\ \frac{d\rho}{dr} = 0 & \Rightarrow \psi = 0 \\ & \frac{d\psi}{d\xi} = 0 \end{aligned}$$

Transformed equation:

$$\frac{1}{r^2} \frac{d}{dr} \left(\frac{r^2}{\rho} \frac{d\rho}{dr} \right) = - \frac{4\pi G m \rho}{kT}$$

$$\begin{aligned} \cancel{\beta^{-1}} \cancel{\lambda} \xi^{-2} \cancel{\beta^{-1/2}} \cancel{\lambda^{1/2}} \frac{d}{d\xi} \left(\cancel{\beta} \cancel{\lambda} \xi^{-1} \cancel{\lambda}^{-1} \exp(+\psi) \right) \cancel{\beta^{-1/2}} \cancel{\lambda}^{1/2} \frac{d}{d\xi} \left(\lambda \exp(-\psi) \right) = \\ = - \left(\frac{4\pi G m}{kT} \right) \lambda \exp(-\psi) \end{aligned}$$

$$\xi^{-2} \frac{d}{d\xi} \left[\cancel{\exp(+\psi)} \cancel{\exp(-\psi)} \left(+ \frac{d\psi}{d\xi} \right) \right] = + \exp(-\psi)$$

$$\xi^{-2} \frac{d}{d\xi} \left(\xi^2 \frac{d\psi}{d\xi} \right) = \exp(-\psi)$$

$\lambda \equiv \rho_c$... parameter of the family of isothermal spheres

4. Construction of the p-V curve

- spherical container with volume V_0
- fixed number of particles N_0
- pressure at the boundary p_0
- constant temperature T

- varying $p_0 \rightarrow$ variation of V_0 and λ
- impossible to find formula $p_0(V_0, N_0, T)$
- however, $\left(\frac{\partial p_0}{\partial V_0}\right)_{N_0, T}$ can be found

$$N = \frac{4\pi}{m} \int_0^r \rho g^2 dy \quad \leftarrow, \quad -\frac{dp}{dr} = \frac{4\pi G \rho}{r^2} \int_0^r \rho g^2 dy$$

$$N = \frac{4\pi}{m} \left(-\frac{dp}{dr}\right) \frac{r^2}{4\pi G \rho} = -\frac{4\pi}{m} \frac{kT}{m} \frac{r^2}{4\pi G \rho} \frac{dp}{dr}$$

$$= -\frac{kT}{G m^2} \beta \lambda^{-1} \int_0^{\xi} \beta^{-1/2} \lambda^{1/2} \frac{d}{d\xi} (\lambda \exp(-\psi)) \lambda^{-1} \exp(\psi) \Rightarrow$$

$$N = \frac{4\pi}{m} \beta^{3/2} \lambda^{-1/2} \int_0^{\xi} \frac{d\psi}{d\xi}$$

(2.9) \leftarrow hydrost. equil. included in this eq.

$$\delta N_0 = \frac{4\pi}{m} \beta^{3/2} \underbrace{\delta \left\{ \lambda^{-1/2} \int_0^{\xi} \left(\frac{d\psi}{d\xi}\right)_0 \right\}}_{=0} = 0$$

- const. numb. of particles
- suffix 0 means "at the boundary"

$$\Rightarrow -\frac{1}{2} \lambda^{-3/2} \int_0^{\xi} \left(\frac{d\psi}{d\xi}\right)_0 \delta \lambda + \lambda^{-1/2} \left\{ \frac{d}{d\xi} \left(\int_0^{\xi} \frac{d\psi}{d\xi} \right) \right\}_0 \delta \xi = 0$$

$$= \int_0^{\xi} \exp(-\psi)$$

Transform: $\lambda, \xi \rightarrow \lambda, r_0$

$$\delta \xi = \left(\frac{\partial \xi_0}{\partial r_0}\right)_\lambda \delta r_0 + \left(\frac{\partial \xi_0}{\partial \lambda}\right)_{r_0} \delta \lambda$$

$$\xi_0 = \beta^{-1/2} \lambda^{1/2} r_0 \quad (2.6)$$

$$\Rightarrow \frac{\partial \xi_0}{\partial r_0} = \beta^{-1/2} \lambda^{1/2}, \quad \frac{\partial \xi_0}{\partial \lambda} = \frac{1}{2} \beta^{-1/2} \lambda^{-1/2} r_0 = \frac{1}{2} \lambda^{-1} \xi_0$$

$$-\frac{1}{2} \lambda^{-3/2} \xi_0^2 \left(\frac{d\psi}{d\xi}\right)_0 \delta \lambda + \lambda^{-1/2} \xi_0^2 \exp(-\psi_0) \delta \xi_0 = 0$$

$$\delta \lambda \left[-\frac{1}{2} \lambda^{-3/2} \xi_0^2 \left(\frac{d\psi}{d\xi}\right)_0 + \frac{1}{2} \lambda^{-3/2} \xi_0^3 \exp(-\psi_0) \right] + \beta^{-1/2} \xi_0^2 \exp(-\psi_0) \delta r_0 = 0$$

$$\delta \lambda = \frac{-\beta^{-1/2} \xi_0^2 \exp(-\psi_0) \delta r_0}{-\frac{1}{2} \lambda^{-3/2} \xi_0^2 \left(\frac{d\psi}{d\xi}\right)_0 + \frac{1}{2} \lambda^{-3/2} \xi_0^3 \exp(-\psi_0)}$$

$$\xi_0, \psi_0, \beta \rightarrow r_0, \rho_0, N_0$$

$$\xi_0 = \beta^{-1/2} \lambda^{1/2} r_0$$

$$\exp(-\psi_0) = \rho_0 \lambda^{-1}$$

$$N_0 = \frac{4\pi}{m} \beta^{3/2} \lambda^{-1/2} \xi_0^2 \left(\frac{d\psi}{d\xi}\right)_0 \quad (2.9)$$

$$\delta \lambda = \frac{-\beta^{-1/2} \beta^{-1} \lambda^{1/2} r_0^2 \rho_0 \lambda^{-1} \delta r_0}{-\frac{1}{2} \lambda^{-3/2} N_0 \frac{m}{4\pi} \beta^{-3/2} \lambda^{1/2} + \frac{1}{2} \lambda^{-3/2} \beta^{-3/2} \lambda^{3/2} r_0^3 \rho_0 \lambda^{-1}}$$

$$= \frac{-2\lambda^2 \rho_0 \delta r_0}{-\frac{m N_0}{4\pi} + r_0^3 \rho_0}$$

$$\delta \lambda = \frac{2\lambda^2 \rho_0 \delta r_0}{\frac{m N_0}{4\pi r_0^2} + r_0^3 \rho_0}$$

$$\left(\frac{\partial \rho_0}{\partial V_0}\right)_{N_0, T} = \frac{kT}{m} \left(\frac{\partial \rho_0}{\partial V_0}\right)_{N_0, T} = \frac{kT}{4\pi m r_0^2} \left(\frac{\partial \rho_0}{\partial r_0}\right)_{N_0, T}$$

$$V = \frac{4\pi}{3} r^3 \rightarrow \frac{\partial V}{\partial r} = 4\pi r^2$$

$$\delta \rho_0 = \delta(\lambda \exp(-\psi_0)) = \exp(-\psi_0) \delta \lambda - \lambda \exp(-\psi_0) \left(\frac{d\psi}{dF}\right)_0 \delta F_0 =$$

~~$\delta \lambda \exp(-\psi_0) - \lambda \exp(-\psi_0) \frac{d\psi}{dF} \delta F_0$~~

$$= \exp(-\psi_0) \left\{ \delta \lambda \left[1 - \lambda \left(\frac{d\psi}{dF}\right)_0 \frac{\partial F_0}{\partial \lambda} \right] - \lambda \left(\frac{d\psi}{dF}\right)_0 \frac{\partial F_0}{\partial r_0} \delta r_0 \right\}$$

$$\left(\frac{d\psi}{dF}\right)_0 = \frac{m N_0}{4\pi} \beta^{-1/2} \lambda^{1/2} \xi_0^{-2} = \frac{m N_0}{4\pi r_0^2} \beta^{-1/2} \lambda^{-1/2}$$

$$\delta \rho_0 = \exp(-\psi_0) \left\{ \delta \lambda \left[1 - \frac{m N_0}{8\pi r_0^2} \beta^{-1/2} \lambda^{-1/2} \xi_0 \right] - \lambda \frac{m N_0}{4\pi r_0^2} \beta^{-1/2} \lambda^{-1/2} \beta^{-1/2} \lambda^{1/2} \delta r_0 \right\}$$

$$\delta \rho_0 = \exp(-\psi_0) \left\{ \delta \lambda \left[1 - \frac{m N_0}{8\pi \beta r_0} \right] - \frac{m N_0 \lambda}{4\pi r_0^2 \beta} \delta r_0 \right\}$$

$$\rho_0 \lambda^{-1} \quad \delta \lambda = \frac{2\lambda \rho_0}{\frac{m N_0}{4\pi r_0^2} - r_0 \rho_0} \delta r_0$$

$$\delta \rho_0 = \rho_0 \lambda^{-1} \left\{ \frac{2\lambda \rho_0}{\frac{m N_0}{4\pi r_0^2} - r_0 \rho_0} \left[1 - \frac{m N_0}{8\pi \beta r_0} \right] - \frac{m N_0 \lambda}{4\pi r_0^2 \beta} \right\} \delta r_0$$

$$\rho_0 = \frac{\rho_0 kT}{m}, \quad \beta = \frac{kT}{4\pi G m}, \quad r_0 = \left[\frac{3}{4\pi} V_0 \right]^{1/3}$$

$$\left(\frac{\partial P_0}{\partial V_0}\right)_{N_0, T} = \frac{kT}{4\pi m r_0^2} \frac{P_0}{r_0} \frac{r_0}{P_0} \left(\frac{\partial P_0}{\partial r_0}\right)_{N_0, T} = \frac{P_0}{3V_0} \frac{r_0}{P_0} \left(\frac{\partial P_0}{\partial r_0}\right)_{N_0, T} =$$

$$= \frac{2P_0}{3V_0} \left\{ \frac{P_0 r_0}{\frac{mN_0}{4\pi r_0^2} - r_0 P_0} \left[1 - \frac{mN_0}{8\pi \beta V_0} \right] - \frac{mN_0}{8\pi r_0^2 \beta} \right\} =$$

$$= \frac{2P_0}{3V_0} \left\{ \frac{r_0 P_0 \left[1 - \frac{Gm^2 N_0}{2kT r_0} \right]}{\frac{mN_0 - 4\pi r_0^3 P_0}{4\pi r_0^2}} - \frac{mN_0 \cdot 4\pi Gm}{2 \cdot 8\pi r_0^2 kT} \right\} =$$

$$= \frac{2P_0}{3V_0} \left\{ \frac{4\pi r_0^3 P_0 (2kT r_0 - Gm^2 N_0)}{2kT r_0 (mN_0 - 4\pi r_0^3 P_0)} - \frac{Gm^2 N_0}{2r_0^2 kT} \right\} =$$

$$= \frac{2P_0}{3V_0} \frac{8\pi r_0^4 P_0 kT - 4\pi r_0^3 P_0 Gm^2 N_0 - Gm^3 N_0^2 + 4\pi r_0^3 P_0 Gm^2 N_0}{2kT r_0 (mN_0 - 4\pi r_0^3 P_0)} =$$

$$= - \frac{2P_0}{3V_0} \frac{Gm^3 N_0^2 - 8\pi r_0^4 P_0 kT}{2kT r_0 mN_0 - 8\pi r_0^4 P_0 kT} = \quad 4\pi r_0^4 = 3V_0 \left(\frac{3V_0}{4\pi}\right)^{1/3}$$

$$= - \frac{2P_0}{3V_0} \frac{1 - \frac{Gm^3 N_0^2}{8\pi r_0^4 P_0 kT}}{1 - \frac{2kT r_0 mN_0}{4\pi r_0^4 P_0 kT}} = - \frac{2P_0}{3V_0} \frac{1 - \frac{Gm^2 N_0^2}{8\pi r_0^4 P_0}}{1 - \frac{N_0 kT}{(4\pi r_0^3 P_0) \cdot 3V_0}} \Rightarrow$$

$$\left(\frac{\partial P_0}{\partial V_0}\right)_{N_0, T} = - \frac{2P_0}{3V_0} \frac{1 - \left(\frac{4\pi}{3}\right)^{1/3} \frac{Gm^2 N_0^2}{6P_0 V_0^{2/3}}}{1 - \frac{N_0 kT}{3P_0 V_0}}$$

for $G=0$: $\left(\frac{\partial P_0}{\partial V_0}\right)_{N_0, T} = - \frac{2P_0}{3V_0} \frac{1}{\frac{3P_0 V_0 - N_0 kT}{3P_0 V_0}} = - \frac{2N_0 kT}{3V_0^2} \frac{3N_0 kT}{2N_0 kT} = - \frac{N_0 kT}{V_0^2}$

(8)

5. Obtaining ρ and V from Emden's eq. solution- ρ and V as functions of ξ , ψ and $\frac{d\psi}{d\xi}$

$$N = \frac{4\pi}{m} \beta^{3/2} \lambda^{-1/2} \xi^2 \frac{d\psi}{d\xi}$$

$$\lambda^{1/2} = \frac{4\pi}{mN} \beta^{3/2} \xi^2 \frac{d\psi}{d\xi}$$

$$\rho = \frac{\rho kT}{m} = \lambda \exp(-\psi) \frac{kT}{m} = \frac{16\pi^2}{m^2 N^2} \left(\frac{kT}{4\pi G m} \right)^3 \xi^4 \left(\frac{d\psi}{d\xi} \right)^2 \exp(-\psi) \frac{kT}{m}$$

$$\rho = \frac{k^4 T^4}{4\pi G^3 m^6 N^2} \xi^4 \left(\frac{d\psi}{d\xi} \right)^2 \exp(-\psi)$$

$$V = \frac{4\pi}{3} r^3 = \frac{4\pi}{3} \beta^{3/2} \lambda^{-3/2} \xi^3 = \frac{4\pi}{3} \beta^{3/2} \left(\frac{mN}{4\pi} \right)^3 \beta^{-9/2} \xi^{-6} \left(\frac{d\psi}{d\xi} \right)^{-3} \xi^3$$

$$= \frac{4\pi}{3} \left(\frac{4\pi G m}{kT} \right)^3 \left(\frac{mN}{4\pi} \right)^3 \xi^{-3} \left(\frac{d\psi}{d\xi} \right)^{-3}$$

$$V = \frac{4\pi}{3} \left(\frac{G m^2 N}{kT \xi \frac{d\psi}{d\xi}} \right)^3$$

6. Singular point of eq. for $\frac{\partial p_0}{\partial V_0}$

$$1 - \left(\frac{4\pi}{3} \right)^{1/3} \frac{G m^2 N_0^2}{G p_0 V_0^{4/3}} = 0$$

$$1 - \frac{N_0 kT}{3 p_0 V_0} = 0 \quad \rightarrow p_0 = \frac{N_0 kT}{3 V_0}$$

$$1 = \left(\frac{4\pi}{3} \right)^{1/3} \frac{G m^2 N_0^2 \cdot 3 V_0}{2 G N_0 kT V_0^{4/3}}$$

$$V_0 = \frac{\pi}{6} \left(\frac{G m^2 N_0}{kT} \right)^3$$

$$p_0 = \frac{2}{\pi} \frac{k^4 T^4}{G^3 m^6 N_0^2}$$

7. Stability

(9)

$$1 - \frac{\left(\frac{4\pi}{3}\right)^{1/2} G m^2 N^2}{6 p V^{4/3}} = 1 - \frac{\left(\frac{4\pi}{3}\right)^{1/2} G m^2 N^2}{6 \frac{k T^4}{4\pi G^3 m^6 N^2} \xi^4 \left(\frac{d\psi}{d\xi}\right)^2 \exp(-\psi) \left(\frac{4\pi}{3}\right)^{1/2} \left(\frac{G m^2 N}{k T \xi \frac{d\psi}{d\xi}}\right)^4}$$

$$= 1 - \frac{G m^2 N^2 \left(\frac{d\psi}{d\xi}\right)^2 \exp(\psi)}{2 G m^2 N^2} = 1 - \frac{1}{2} \exp(\psi) \left(\frac{d\psi}{d\xi}\right)^2 \quad (3.1)$$

$$1 - \frac{k T N}{3 p V} = 1 - \frac{k T N}{3 \frac{k T^4}{4\pi G^3 m^6 N^2} \xi^4 \left(\frac{d\psi}{d\xi}\right)^2 \exp(-\psi) \left(\frac{4\pi}{3}\right)^{1/2} \left(\frac{G m^2 N}{k T \xi \frac{d\psi}{d\xi}}\right)^3} =$$

$$= 1 - \xi^{-1} \exp(\psi) \frac{d\psi}{d\xi} \quad (3.2)$$

$$\left(\frac{\partial p}{\partial V}\right)_{N,T} \equiv \frac{(3.1)}{(3.2)}$$

both (3.1) and (3.2) are positive for $\xi < 6,5$

critical radius: $r_{\text{crit}} = 6,5 \beta^{1/2} \lambda^{-1/2} = 6,5 \left(\frac{k T}{4\pi G m \lambda}\right)^{1/2}$

$$r_{\text{crit}} = 1,8 \left(\frac{k T}{m G \rho_{\text{cen}}}\right)^{1/2}$$

integrating p :

$$\frac{M}{r_{\text{crit}}} = 2,4 \frac{k T}{m G}$$

compare to Jeans relation:

$$\frac{M}{r_{\text{Jeans}}} = 3,3 \frac{k T}{m G}$$

Summary:

Emden eq.: $\frac{1}{r^2} \frac{d}{dr} \left(\frac{r^2}{\rho} \frac{d\rho}{dr} \right) = - \frac{4\pi G m \rho}{kT}$

Transformation: $r, \rho, T, m \rightarrow \xi, \psi, \lambda, \beta$

$r = \beta^{1/2} \lambda^{-1/2} \xi, \rho = \lambda \exp(-\psi), \beta = \frac{kT}{4\pi G m}$

$\rightarrow \xi^{-2} \frac{d}{d\xi} \left(\xi^2 \frac{d\psi}{d\xi} \right) = \exp(-\psi)$

+ boundary cond at $r = \phi$: $\lambda = \rho_c, \frac{d\psi}{d\xi} = 0, \psi = \phi$

Variations at the boundary for fixed N_0 :

$N_0 = \frac{4\pi}{m} \beta^{3/2} \lambda^{-1/2} \xi^2 \left(\frac{d\psi}{d\xi} \right)_0$ - From hydrost. equil.

$\delta N_0 \left(\lambda, \xi_0, \left(\frac{d\psi}{d\xi} \right)_0 \right) = 0$

$\lambda^{1/2} = \frac{4\pi}{mN} \beta^{3/2} \xi^2 \frac{d\psi}{d\xi}$
 - to determine $\lambda \equiv \rho$
 - ξ and ψ obtained from P-V diagram

$\square \delta \lambda + \square \delta \xi_0 = 0 \xrightarrow{\xi_0 \rightarrow r_0} \square \delta \lambda + \square \delta r_0 = 0$

$\rightarrow \delta \lambda = \square \delta r_0$

$\xi_0, \psi_0, \beta \rightarrow r_0, \rho_0, N_0$

P-V curve: $\left(\frac{\partial \rho_0}{\partial V_0} \right)_{N_0, T} = \frac{kT}{m} \left(\frac{\partial \rho_0}{\partial V_0} \right)_{N_0, T} = \frac{kT}{4\pi m r_0^2} \left(\frac{\partial \rho_0}{\partial r_0} \right)_{N_0, T}$

$\delta \rho_0 = \delta (\lambda \exp(-\psi_0)) = \square \delta \lambda + \square \delta \xi_0 \xrightarrow{\xi_0 \rightarrow r_0}$

$= \square \delta \lambda + \square \delta r_0 \rightarrow \delta \rho_0 = \square \delta r_0$

$\left(\frac{\partial \rho_0}{\partial V_0} \right)_{N_0, T} = - \frac{2\rho_0}{3V_0} \frac{1 - \left(\frac{4\pi T}{3} \right)^{1/3} \frac{G m^2 N_0^2}{6\rho_0 V_0}}{1 - \frac{N_0 k T}{3\rho_0 V_0}}$